

Quantum Mechanics 742 – Second Test
Covering Chapters 14 – 18

The following new equations you should memorize, and understand how to use them:

Cross-section $\Gamma = n\sigma \Delta \mathbf{v} $ $\sigma = \int \frac{d\sigma}{d\Omega} d\Omega$	Momentum Change $\mathbf{K} = k\hat{\mathbf{r}} - k\hat{\mathbf{z}}$	Fermi's Golden Rule: $\mathcal{T}_{FI} = W_{FI} + \dots$ $\Gamma(I \rightarrow F) = 2\pi\hbar^{-1} \mathcal{T}_{FI} ^2 \delta(E_F - E_I)$	Time Dependent Perturbation Theory $H = H_0 + W(t)$ $H_0 \phi_n\rangle = E_n \phi_n\rangle$ $\omega_{nm} \equiv (E_n - E_m)/\hbar$ $W_{nm}(t) \equiv \langle \phi_n W(t) \phi_m \rangle$ $P(I \rightarrow F) = S_{FI} ^2$
Coulomb Gauge $\nabla \cdot \mathbf{A} = 0$	Sudden: $T\Delta E \ll \hbar$ $P(I \rightarrow F) = \langle \psi_F \psi_I \rangle ^2$ Adiabatic: $T\Delta E \gg \hbar$ $P(I \rightarrow F) = \delta_{IF}$	Photon Creation and Annihilation Operators $[a_{\mathbf{k}\sigma}, a_{\mathbf{k}'\sigma'}^\dagger] = \delta_{\mathbf{k},\mathbf{k}'} \delta_{\sigma,\sigma'}$ $a_{\mathbf{k}\sigma} n, \mathbf{k}, \sigma\rangle = \sqrt{n} n-1, \mathbf{k}, \sigma\rangle$ $a_{\mathbf{k}\sigma}^\dagger n, \mathbf{k}, \sigma\rangle = \sqrt{n+1} n+1, \mathbf{k}, \sigma\rangle$	EM Hamiltonian $H = \sum_{\mathbf{k},\sigma} \hbar\omega_k a_{\mathbf{k}\sigma}^\dagger a_{\mathbf{k}\sigma}$
Electric Dipole Moment $\mathbf{r}_{nm} \equiv \langle \phi_n \mathbf{R} \phi_m \rangle$	EM waves $\boldsymbol{\epsilon}_{\mathbf{k}\sigma} \cdot \boldsymbol{\epsilon}_{\mathbf{k}\sigma'}^* = \delta_{\sigma\sigma'}$ $\boldsymbol{\epsilon}_{\mathbf{k}\sigma} \cdot \mathbf{k} = 0$ $\omega_k = ck$	Converting Finite \rightarrow Infinite $\lim_{V \rightarrow \infty} \left[\frac{1}{V} \sum_{\mathbf{k}} f(\mathbf{k}) \right] = \int f(\mathbf{k}) \frac{d^3\mathbf{k}}{(2\pi)^3}$	

Other things you should know:

- The meaning of cross-section and differential cross-section
- Doing integrals in spherical coordinates
- How to do integrals like $\int f(x) \delta[g(x)] dx$
- In the adiabatic approximation, make sure the eigenstates you use are correct eigenstates of the initial and final Hamiltonian.
- In the adiabatic approximation: The lowest energy state goes to the lowest state, second lowest to second lowest, etc.
 - However, if there is a symmetry that always commutes with the Hamiltonian, then first sort states by eigenvalues of that symmetry
- For harmonic perturbations, make sure you know how to extract W given $W(t)$, and don't get the two confused
 - In contrast, when the perturbation is independent of time, $W(t) = W$
- How to compute expressions like $\boldsymbol{\epsilon}_{\mathbf{k}\sigma} \cdot \mathbf{r}_{FI}$ for each of the two possible polarizations
- How to average (or sum) over polarizations, and average (or integrate) over angles
- How to take the limit $V \rightarrow \infty$, and turn sums over final states into integrals you can do
- What a quantum state like $|n_1, \mathbf{k}_1, \sigma_1; n_2, \mathbf{k}_2, \sigma_2\rangle$ means, or $|\phi_a; n_1, \mathbf{k}_1, \sigma_1; n_2, \mathbf{k}_2, \sigma_2\rangle$ means
- Understanding that after quantizing the EM field, electric and magnetic fields are now operators that are functions of \mathbf{r} , but not t
- How to write out sums with EM fields and expectation values where the sums collapse to few terms
- How to find the energy of a system of photons, or photons plus an atom
 - Related: the time dependence of such a system
- How to create or annihilate one photon from any state with any number of photons
- Qualitatively, what our diagrams mean in our computations

- Why we concentrate only on certain diagrams when scattering near a resonance

The following new equations you need not memorize, but you should know how to use them if given to you:

Time-dependent Perturbation Theory			
$S_{FI} = \delta_{FI} + (i\hbar)^{-1} \int_0^T dt W_{FI}(t) e^{i\omega_{FI}t} + (i\hbar)^{-2} \sum_n \int_0^T dt W_{Fn}(t) e^{i\omega_{Fn}t} \int_0^t dt' W_{nl}(t') e^{i\omega_{nl}t'} + \dots$			
Electric Dipole Absorption $\Gamma_{EI} = 4\pi^2 \alpha \hbar^{-1} \mathcal{I}(\omega_{FI}) \mathbf{\epsilon} \cdot \mathbf{r}_{FI} ^2$		Harmonic Perturbations: $W(t) = W e^{-i\omega t} + W^\dagger e^{i\omega t}, \quad W_{FI} \equiv \langle \phi_F W \phi_I \rangle$	
Scattering Away from Resonance $\mathcal{T}_{FI} = \frac{-e^2}{\epsilon_0 V} \sum_n \frac{\omega \omega_{nl}}{\omega_n^2 - \omega^2} (\mathbf{\epsilon}_F^* \cdot \mathbf{r}_{nl}^*) (\mathbf{\epsilon}_I \cdot \mathbf{r}_{nl})$		$\Gamma(I \rightarrow F) = \begin{cases} 2\pi \hbar^{-1} W_{FI} ^2 \delta(E_F - E_I - \hbar\omega) & \text{if } E_F > E_I \\ 2\pi \hbar^{-1} W_{FI}^\dagger ^2 \delta(E_F - E_I + \hbar\omega) & \text{if } E_F < E_I \end{cases}$	
Spontaneous Decay $\Gamma = \frac{4\alpha}{3c^2} \omega_{IF}^3 \mathbf{r}_{FI} ^2$	Transition matrix: $\mathcal{T}_{FI} = W_{FI} + \lim_{\epsilon \rightarrow 0^+} \left[\sum_n \frac{W_{Fn} W_{nl}}{(E_I - E_n + i\epsilon)} + \sum_n \sum_m \frac{W_{Fn} W_{nm} W_{ml}}{(E_I - E_m + i\epsilon)(E_I - E_n + i\epsilon)} + \dots \right]$		
Dirac Equation $i\hbar \frac{\partial}{\partial t} \Psi = c\boldsymbol{\alpha} \cdot (\mathbf{P} + e\mathbf{A}) \Psi - eU\Psi + mc^2 \beta \Psi$ $\beta = \begin{pmatrix} \mathbf{0} & -\mathbf{1} \\ -\mathbf{1} & \mathbf{0} \end{pmatrix}, \quad \boldsymbol{\alpha} = \begin{pmatrix} \boldsymbol{\sigma} & \mathbf{0} \\ \mathbf{0} & -\boldsymbol{\sigma} \end{pmatrix}$ $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$		Electromagnetic Operators $\mathbf{A}(\mathbf{r}) = \sum_{\mathbf{k}, \sigma} \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} (a_{\mathbf{k}\sigma} \boldsymbol{\epsilon}_{\mathbf{k}\sigma} e^{i\mathbf{k}\cdot\mathbf{r}} + a_{\mathbf{k}\sigma}^\dagger \boldsymbol{\epsilon}_{\mathbf{k}\sigma}^* e^{-i\mathbf{k}\cdot\mathbf{r}})$ $\mathbf{E}(\mathbf{r}) = \sum_{\mathbf{k}, \sigma} \sqrt{\frac{\hbar \omega_k}{2\epsilon_0 V}} i (a_{\mathbf{k}\sigma} \boldsymbol{\epsilon}_{\mathbf{k}\sigma} e^{i\mathbf{k}\cdot\mathbf{r}} - a_{\mathbf{k}\sigma}^\dagger \boldsymbol{\epsilon}_{\mathbf{k}\sigma}^* e^{-i\mathbf{k}\cdot\mathbf{r}})$ $\mathbf{B}(\mathbf{r}) = \sum_{\mathbf{k}, \sigma} \sqrt{\frac{\hbar}{2\epsilon_0 V \omega_k}} i \mathbf{k} \times (a_{\mathbf{k}\sigma} \boldsymbol{\epsilon}_{\mathbf{k}\sigma} e^{i\mathbf{k}\cdot\mathbf{r}} - a_{\mathbf{k}\sigma}^\dagger \boldsymbol{\epsilon}_{\mathbf{k}\sigma}^* e^{-i\mathbf{k}\cdot\mathbf{r}})$	
Lab vs. CM frame $\left(\frac{d\sigma}{d\Omega} \right)_L = \left(\frac{d\sigma}{d\Omega} \right) \frac{(1 + 2\gamma \cos \theta + \gamma^2)^{3/2}}{1 + \gamma \cos \theta}$ $\gamma = \frac{m}{M}$		Scattering Formulas $k^2 = \frac{2\mu E}{\hbar^2}$	Scattering Near Resonance $\sigma = \frac{8\pi \alpha^2 \omega_{nl}^4 \mathbf{r}_{nl} ^4}{3c^2 \left[(\omega - \omega_{nl})^2 + \frac{1}{4} \Gamma^2 \right]}$
1 st Born Approximation $\frac{d\sigma}{d\Omega} = \frac{\mu^2}{4\pi^2 \hbar^4} \left \int d^3 \mathbf{r} V(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{r}} \right ^2$			