## Physics 741 - Graduate Quantum Mechanics 1

## Solutions to Chapter 4

4.2 [10] The commutation relations of the angular momentum operator $L_{z}$ with the momentum operator $P$ were to be worked out in chapter 3, problem 3.
(a) [3] Using these commutation relations, derive two non-trivial uncertainty relationships.

The commutation relations were

$$
\left[L_{z}, P_{x}\right]=i \hbar P_{y}, \quad\left[L_{z}, P_{y}\right]=-i \hbar P_{x}, \quad\left[L_{z}, P_{z}\right]=0
$$

According to the generalized uncertainty relation, we therefore have

$$
\Delta L_{z} \Delta P_{x} \geq \frac{1}{2} \hbar\left|\left\langle P_{y}\right\rangle\right|, \quad \Delta L_{z} \Delta P_{y} \geq \frac{1}{2} \hbar\left|\left\langle P_{x}\right\rangle\right|, \quad \Delta L_{z} \Delta P_{z} \geq 0
$$

Since the left side of each of these expressions is the product of two positive numbers, the last inequality doesn't give us any information, but the other two inequalities suffice.
(b) [3] Show that if you are in an eigenstate of any observable, the uncertainty in that observable is zero.

Let $A$ be any observable, and $|a\rangle$ a normalized eigenstate with $A|a\rangle=a|a\rangle$. Then
$(\Delta A)^{2}=\left\langle A^{2}\right\rangle-\langle A\rangle^{2}=\langle a| A^{2}|a\rangle-\langle a| A|a\rangle^{2}=a\langle a| A|a\rangle-(a\langle a \mid a\rangle)^{2}=a^{2}\langle a \mid a\rangle-a^{2}=0$
(c) [4] Show that if you are in an eigenstate of $\boldsymbol{L}_{\boldsymbol{x}}$, then you must have $\left\langle P_{x}\right\rangle=\left\langle P_{y}\right\rangle=0$.

According to our inequalities we found above, if we are in an eigenstate of $L_{z}$, then we have $\Delta L_{z}=0$, and therefore

$$
0 \geq \frac{1}{2} \hbar\left|\left\langle P_{y}\right\rangle\right|, \quad 0 \geq \frac{1}{2} \hbar\left|\left\langle P_{x}\right\rangle\right| .
$$

Since absolute values are never negative, it follows that $\left|\left\langle P_{y}\right\rangle\right|=\left|\left\langle P_{x}\right\rangle\right|=0$, which implies

$$
\left\langle P_{x}\right\rangle=\left\langle P_{y}\right\rangle=0 .
$$

4.3 [10] We will eventually discover that particles have spin, which is described by three operators $\mathbf{S}=\left(S_{x}, S_{y}, S_{z}\right)$ with commutation relations

$$
\left[S_{x}, S_{y}\right]=i \hbar S_{z}, \quad\left[S_{y}, S_{z}\right]=i \hbar S_{x}, \quad\left[S_{z}, S_{x}\right]=i \hbar S_{y}
$$

A particle in a magnetic field of magnitude $B$ pointing in the $\boldsymbol{z}$-direction will have Hamiltonian $H=-\mu B S_{z}$ where $\mu$ and $B$ are constants.
(a) [5] Derive formulas for the time derivative of all three components of $\langle\mathbf{S}\rangle$

We use the standard formula for the time evolution of any operator, namely

$$
\begin{aligned}
& \frac{d}{d t}\left\langle S_{x}\right\rangle=\frac{i}{\hbar}\left\langle\left[H, S_{x}\right]\right\rangle=-\frac{i \mu B}{\hbar}\left\langle\left[S_{z}, S_{x}\right]\right\rangle=-\frac{i \mu B}{\hbar} i \hbar\left\langle S_{y}\right\rangle=\mu B\left\langle S_{y}\right\rangle, \\
& \frac{d}{d t}\left\langle S_{y}\right\rangle=\frac{i}{\hbar}\left\langle\left[H, S_{y}\right]\right\rangle=-\frac{i \mu B}{\hbar}\left\langle\left[S_{z}, S_{y}\right]\right\rangle=-\frac{i \mu B}{\hbar}(-i \hbar)\left\langle S_{x}\right\rangle=-\mu B\left\langle S_{x}\right\rangle, \\
& \frac{d}{d t}\left\langle S_{z}\right\rangle=\frac{i}{\hbar}\left\langle\left[H, S_{z}\right]\right\rangle=-\frac{i \mu B}{\hbar}\left\langle\left[S_{z}, S_{z}\right]\right\rangle=-\frac{i \mu B}{\hbar} 0=0 .
\end{aligned}
$$

(b) [5] At time $t=0$, the expectation values of $\langle S\rangle$ are given by

$$
\left\langle S_{x}\right\rangle_{t=0}=a, \quad\left\langle S_{y}\right\rangle_{t=0}=0, \quad\left\langle S_{z}\right\rangle_{t=0}=b
$$

Determine the expectation value $\langle\mathbf{S}\rangle_{t}$ at later times.
Since the time derivative of $\left\langle S_{z}\right\rangle$ vanishes, this will just remain constant. The expectation value of the other two operators, however, are related by

$$
d\left\langle S_{x}\right\rangle / d t=\mu B\left\langle S_{y}\right\rangle, \quad d\left\langle S_{y}\right\rangle / d t=-\mu B\left\langle S_{x}\right\rangle .
$$

One way to proceed is to take another derivative of the first equation, which yields

$$
d^{2}\left\langle S_{x}\right\rangle / d t^{2}=\mu B d\left\langle S_{y}\right\rangle / d t=-\mu^{2} B^{2}\left\langle S_{x}\right\rangle .
$$

This suggests solutions along the lines of

$$
\left\langle S_{x}\right\rangle=\alpha \cos (\mu B t)+\beta \sin (\mu B t)
$$

Taking the derivative we can find

$$
\left\langle S_{y}\right\rangle=-\alpha \sin (\mu B t)+\beta \cos (\mu B t) .
$$

Our boundary conditions at $t=0$ then tell us that $\alpha=a$ and $\beta=0$. In summary, we have

$$
\left\langle S_{x}\right\rangle_{t}=a \cos (\mu B t), \quad\left\langle S_{y}\right\rangle_{t}=-a \sin (\mu B t), \quad\left\langle S_{z}\right\rangle_{t}=b .
$$

