

# Lecture #7: Method of Multiple Scales

Example:

$$\ddot{y} + \varepsilon \dot{y} + y = 0$$

$$y(0) = 0$$

$$\dot{y}(0) = 1$$

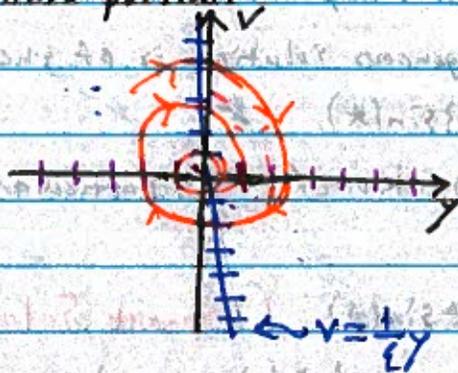
We can draw a phase portrait

$$\dot{y} = v$$

$$\dot{v} = -\varepsilon v - y$$

$$J(0,0) = \begin{bmatrix} 0 & 1 \\ -1 & -\varepsilon \end{bmatrix}$$

⇒ stable spirals



The exact solution is given by

$$y(t) = \frac{1}{\sqrt{1 - \varepsilon^2/4}} e^{-\varepsilon t/2} \sin(t \sqrt{1 - \varepsilon^2/4})$$

period =  $\frac{2\pi}{\sqrt{1 - \varepsilon^2/4}} \sim 2\pi(1 - \varepsilon^2/8 + \dots)$   
 Damping Scaling =  $1/\varepsilon$

Lets try a naive solution

$$y = y_0 + \varepsilon y_1 + \varepsilon^2 y_2 + \dots$$

$O(1)$ :

$$\ddot{y}_0 + y_0 = 0$$

$$y_0(0) = 0, \dot{y}_0(0) = 1$$

$$\Rightarrow y_0(t) = \sin(t)$$

OCEI:

$$\ddot{y}_1 + \gamma_0 y_1 = 0, \quad y_1(0) = 0$$
$$\Rightarrow \ddot{y}_1 + \gamma_1 y_1 = -\cos(t), \quad \dot{y}_1(0) = 0$$

Let  $\mathcal{L} = \frac{d^2}{dt^2} + 1$ , we have

$$\mathcal{L} y_1 = -\cos(t).$$

Therefore,

$$\text{Ker}(\mathcal{L}) = \text{span}\{\cos(t), \sin(t)\}.$$

Consequently, the homogeneous solution is of the form

$$y_1^h = A \cos(t) + B \sin(t).$$

However, since  $-\cos(t) \in \text{Ker}(\mathcal{L})$  the particular solution will be of the form

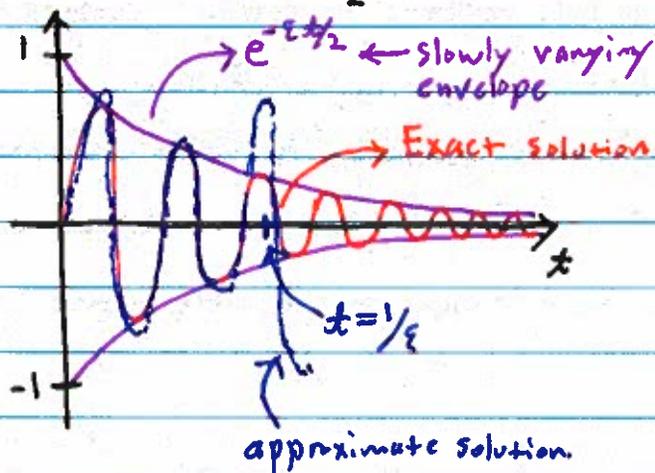
$$y_1^p = C t \cos(t) + D t \sin(t) \quad (\text{Resonant Solutions})$$

Using the initial conditions and solving we have

$$y_1 = y_1^h + y_1^p = \frac{1}{2} t \sin(t)$$

Therefore, to two terms we have that

$$y \sim \sin(t) + \frac{1}{2} \epsilon t \sin(t).$$



The asymptotic expansion works for fixed  $t$  and decreasing  $\epsilon$ . However, what we want is a uniform expansion that works for arbitrarily large times.

The key observation is that there is a second time scale  $\epsilon t$  that slightly distorts the intrinsic period from being resonant. A useful expansion has to account for both of these scales. A better ansatz assumes the solution depends on multiple scales:

$$t_1 = t, t_2 = \epsilon t, t_3 = \epsilon^2 t, \dots$$

$$y(t, \epsilon) = y(t_1, t_2, t_3, \dots, \epsilon)$$

Now,

$$\frac{d}{dt} = \frac{\partial}{\partial t_1} \frac{\partial}{\partial t} + \frac{\partial}{\partial t_2} \frac{\partial}{\partial t} + \frac{\partial}{\partial t_3} \frac{\partial}{\partial t} + \dots$$

$$= \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2} + \epsilon^2 \frac{\partial}{\partial t_3} + \dots$$

$$\frac{d^2}{dt^2} = \left( \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2} + \epsilon^2 \frac{\partial}{\partial t_3} + \dots \right) \left( \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2} + \epsilon^2 \frac{\partial}{\partial t_3} + \dots \right)$$

$$= \frac{\partial^2}{\partial t_1^2} + \epsilon \left( \frac{\partial^2}{\partial t_1 \partial t_2} + \frac{\partial^2}{\partial t_2 \partial t_1} \right) + \epsilon^2 \left( \frac{\partial^3}{\partial t_1^2 \partial t_2} + \frac{\partial^3}{\partial t_1 \partial t_2^2} + \frac{\partial^3}{\partial t_2^2 \partial t_1} \right) + \dots$$

can't always assume mixed partials commute.

but we will and see if things go wrong.

The ODE up to  $\epsilon^2$  becomes

$$\frac{\partial^2 y}{\partial t_1^2} + 2\epsilon \frac{\partial^2 y}{\partial t_1 \partial t_2} + \epsilon^2 \frac{\partial^2 y}{\partial t_2^2} + 2\epsilon^2 \frac{\partial^3 y}{\partial t_1 \partial t_2^2} + \epsilon \frac{\partial y}{\partial t_1} + \epsilon^2 \frac{\partial y}{\partial t_2} + y = 0$$

$$y(0) = 0$$

$$\dot{y}(0) = 1$$

We now expand

$$y = y_0 + \epsilon y_1 + \epsilon^2 y_2 + \dots$$

$\mathcal{O}(1)$ :

$$\frac{\partial^2 y_0}{\partial t_1^2} + y_0 = 0, \quad y_0(0) = 0, \quad ;$$

$$\Rightarrow y_0(t_1, t_2) = A(t_2) \sin(t_1)$$

$\mathcal{O}(\varepsilon)$ :

$$\frac{\partial^2 y_1}{\partial t_1^2} + 2 \frac{\partial^2 y_1}{\partial t_1 \partial t_2} + \frac{\partial y_1}{\partial t_1} + y_1 = 0$$

$$\Rightarrow \frac{\partial^2 y_1}{\partial t_1^2} + 2 \frac{\partial A}{\partial t_2} \cos(t_1) + A(t_2) \cos(t_1) + y_1 = 0$$

$$\Rightarrow \frac{\partial^2 y_1}{\partial t_1^2} + y_1 = -2 \frac{\partial A}{\partial t_2} \cos(t_1) + A(t_2) \cos(t_1)$$

$$\Rightarrow \left( \frac{\partial^2}{\partial t_1^2} + 1 \right) y_1 = -2 \frac{\partial A}{\partial t_2} \cos(t_1) + A(t_2) \cos(t_1)$$

To remove the secular terms we need

$$-2 \frac{\partial A}{\partial t_2} + A(t_2) = 0$$

$$\Rightarrow A(t_2) = e^{-t_2/2}$$

Therefore, to lowest order

$$y(t) \sim e^{-\frac{1}{2}t} \sin(t)$$

Remarks:

We derived an amplitude equation and a wave equation:

$$\frac{\partial A}{\partial t_2} = -\frac{1}{2} A(t_2) \quad \text{and} \quad \frac{\partial^2 y}{\partial t_1^2} + y_0 = 0$$

Example:

$$\ddot{y} + \epsilon \dot{y}^3 + y = 0$$

Arbitrary initial conditions.

Naive Expansion

$$y = y_0 + \epsilon y_1 + \epsilon^2 y_2 + \dots$$

$O(1)$ :

$$\ddot{y}_0 + y_0 = 0$$

$$\Rightarrow y_0 = A e^{it} + A^* e^{-it}$$

Complex amplitude

Complex conjugate

$O(\epsilon)$ :

$$\ddot{y}_1 + \dot{y}_0^3 + y_1 = 0$$

$$\Rightarrow \ddot{y}_1 + y_1 = - (A i e^{it} - A^* i e^{-it})^3$$

$$\Rightarrow \ddot{y}_1 + y_1 = -i (A e^{it} - A^* e^{-it})^3$$

$$\Rightarrow \ddot{y}_1 + y_1 = -i (A^3 e^{3it} - 3A^2 e^{it} + 3A(A^*)^2 e^{-it} - (A^*)^3 e^{-3it})$$

These terms are resonant

Multiple Scale Expansion

$$\text{Let } t_1 = t, t_2 = \epsilon t.$$

$$\Rightarrow \frac{d}{dt} = \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2}$$

$$\Rightarrow \frac{d^2}{dt^2} = \left( \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2} \right) \left( \frac{\partial}{\partial t_1} + \epsilon \frac{\partial}{\partial t_2} \right)$$

$$= \frac{\partial^2}{\partial t_1^2} + 2\epsilon \frac{\partial^2}{\partial t_1 \partial t_2} + \epsilon^2 \frac{\partial^2}{\partial t_2^2}$$

$$\text{Assume } y(t_1, t_2, \epsilon) = y_0(t_1, t_2) + \epsilon y_1(t_1, t_2) + \dots$$

The ODE becomes the following PDE:

$$\frac{\partial^2 y}{\partial t_1^2} + 2\varepsilon \frac{\partial^2 y}{\partial t_1 \partial t_2} + \varepsilon^2 \frac{\partial^2 y}{\partial t_2^2} + \varepsilon \left( \frac{\partial y}{\partial t_1} + \varepsilon \frac{\partial y}{\partial t_2} \right)^3 + y = 0.$$

The order by order equations are given by

$O(1)$ :

$$\frac{\partial^2 y_0}{\partial t_1^2} + y_0 = 0$$

$$\Rightarrow y_0 = A(t_2) e^{it_1} + A^*(t_2) e^{-it_1}$$

$O(\varepsilon)$ :

$$\frac{\partial^2 y_1}{\partial t_1^2} + 2 \frac{\partial^2 y_1}{\partial t_1 \partial t_2} + \left( \frac{\partial y_0}{\partial t_1} \right)^3 + y_1 = 0$$

$$\Rightarrow \frac{\partial^2 y_1}{\partial t_1^2} + y_1 = -2i \left( \frac{\partial A}{\partial t_2} e^{it_1} - \frac{\partial A^*}{\partial t_2} e^{-it_1} \right) + \left( iA e^{it_1} - iA^* e^{-it_1} \right)^3$$

$$\Rightarrow \frac{\partial^2 y_1}{\partial t_1^2} + y_1 = -2i \left( \frac{\partial A}{\partial t_2} e^{it_1} - \frac{\partial A^*}{\partial t_2} e^{-it_1} \right) + iA^3 e^{3it_1} + 3iA^2 A^* e^{it_1} + 3iA(A^*)^2 e^{-it_1} - i(A^*)^3 e^{-3it_1}$$

To remove resonant terms we need

$$2i \frac{\partial A}{\partial t_2} + 3iA^2 A^* = 0$$

$$2i \frac{\partial A^*}{\partial t_2} + 3i(A^*)^2 A = 0$$

These equations are complex conjugates of each other so we really only need:

$$\frac{\partial A}{\partial t_2} = -\frac{3}{2} A^2 A^*$$

How do we solve complex valued ODEs??

$$\text{Let } A = r(t_2) e^{i\theta(t_2)}$$

$$\Rightarrow \frac{dA}{dt_2} = r' e^{i\theta(t_2)} + r i \theta' e^{i\theta} = -\frac{3}{2} r^3 e^{i\theta}$$

$$\Rightarrow r' + i r \theta' = -\frac{3}{2} r^3$$

Comparing real and imaginary parts we have:

$$r' = -\frac{3}{2} r^3, \quad \theta' = 0$$

$$\Rightarrow \int_{r_0}^r \frac{1}{r^3} dr = \int_0^{t_2} -\frac{3}{2} dt_2$$

$$\Rightarrow -\frac{1}{2} r^{-2} + \frac{1}{2} r_0^{-2} = -\frac{3}{2} t_2$$

$$\Rightarrow r^{-2} = 3 t_2 + r_0^{-2}$$

$$\Rightarrow r = \frac{r_0}{\sqrt{1 + 3 r_0^2 t_2}}$$

The second differential equation in  $\theta'$  yields  $\theta = \theta_0$ .

Putting it all together:

$$\begin{aligned} y \sim y_0 &= A(t_2) e^{i t_2} + A^*(t_2) e^{-i t_2} \\ &= \frac{r_0}{\sqrt{1 + 3 r_0^2 t_2}} e^{i \theta_0} e^{i t_2} + \frac{r_0}{\sqrt{1 + 3 r_0^2 t_2}} e^{-i \theta_0} e^{-i t_2} \end{aligned}$$

$$\Rightarrow y \sim y_0 = \frac{r_0}{\sqrt{1 + 3 r_0^2 t_2}} \left( e^{i(\theta_0 + t_2)} + e^{-i(\theta_0 + t_2)} \right)$$

$$\Rightarrow y \sim y_0 = \frac{2 r_0}{\sqrt{1 + 3 r_0^2 t_2}} \cos(\theta_0 + t_2)$$