PHY 711 Classical Mechanics and Mathematical Methods
10-10:50 AM MWF in Olin 103

Notes on Lecture 33:

Chapter 10 in F & W: Surface waves

1. Water waves in a channel
2. Wave-like solutions; wave speed

29	Fri, 11/03/2023	Chap. 9	Ideal gas fluids	#23
31	Mon, 11/06/2023	Chap. 9	Traveling and standing waves in the linear approximation	#24
32	Wed, 11/08/2023	Chap. 9	Non-linear and other wave properties	Topic
33	Fri, 11/10/2023	Chap. 9	Analysis of non-linear waves and shock waves	#25
34	Mon, 11/13/2023	Chap. 10	Surface waves in fluids	#26
35	Wed, 11/15/2023	Chap. 10	Surface waves in fluids; soliton solutions	
36	Fri, 11/17/2023	Chap. 11	Heat conduction	
37	Mon, 11/20/2023	Chap. 12	Viscous effects in hydrodynamics	
	Wed, 11/22/2023	Thanksgiving		
	Fri, 11/24/2023	Thanksgiving		
	Mon, 11/27/2023		Presentations I	
	Wed, 11/29/2023		Presentations 2	
	Fri, 12/01/2023		Presentations 3	
38	Mon, 12/04/2023	Chap. 12	Viscous effects in hydrodynamics	
39	Wed, 12/06/2023		Review	
40	Fri, 12/08/2023		Review	

PHY 711 -- Assignment #26

Assigned: 11/13/2023 Due: 11/20/2023

Start reading Chapter 10 in **Fetter & Walecka**.

1. Work Problem 10.3 at the end of Chapter 10 in **Fetter and Walecka**. Note that some of ideas are discussed in Lecture 33.

Immediately after class you might want to sign up for your class presentation times on the google doc --

PHY 711 Presentation Schedule

Monday 11/27/2023

	Presenter Name	Topic
10:00-10:16		
10:17-10:33		
10:34-10:50		

Wednesday 11/29/2023

	Presenter Name	Topic
10:00-10:16		
10:17-10:33		
10:34-10:50		

Friday 12/01/2023

	Presenter Name	Topic
10:00-10:16		
10:17-10:33		
10:34-10:50		

Reference: Chapter 10 of Fetter and Walecka

Physics of incompressible fluids and their surfaces

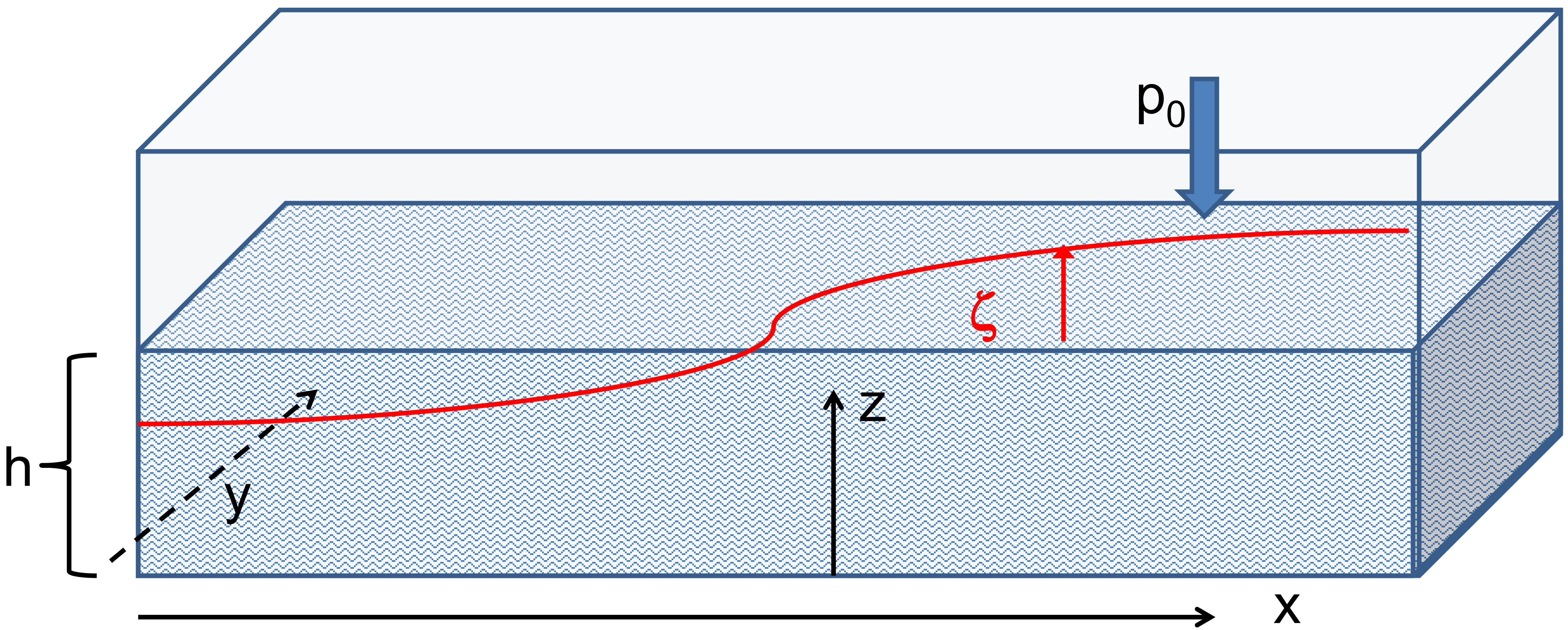
Consider a container of water with average height h and surface h+ζ(x,y,t); (h ⇐⇒ z_0 on some of the slides)

Atmospheric pressure is in equilibrium with the surface of water

Pressure at a height z above the bottom where the surface is at a height $h+\zeta$:

$$p(z) = \begin{cases} p_0 + \rho g\left(h+\zeta-z\right) & \text{For } z \le h+\zeta \\ p_0 & \text{For } z > h+\zeta \end{cases}$$

Here ρ represents density of water



Why do we not consider ρ_{air} in this analysis?

a. Because it is a reasonable approximation
b. Because it simplifies the analysis
c. Both of the above

Euler's equation inside a incompressible fluid:

$$\frac{d\mathbf{v}}{dt} = f_{applied} - \frac{\nabla p}{\rho} = -g\hat{\mathbf{z}} - \frac{\nabla p}{\rho}$$

Assume that $v_z << v_x, v_y$ $\quad \Rightarrow -g - \frac{1}{\rho}\frac{\partial p}{\partial z} \approx 0$

$$\Rightarrow p(x,y,z,t) = p_0 + \rho g\left(\zeta(x,y,t) + h - z\right)$$

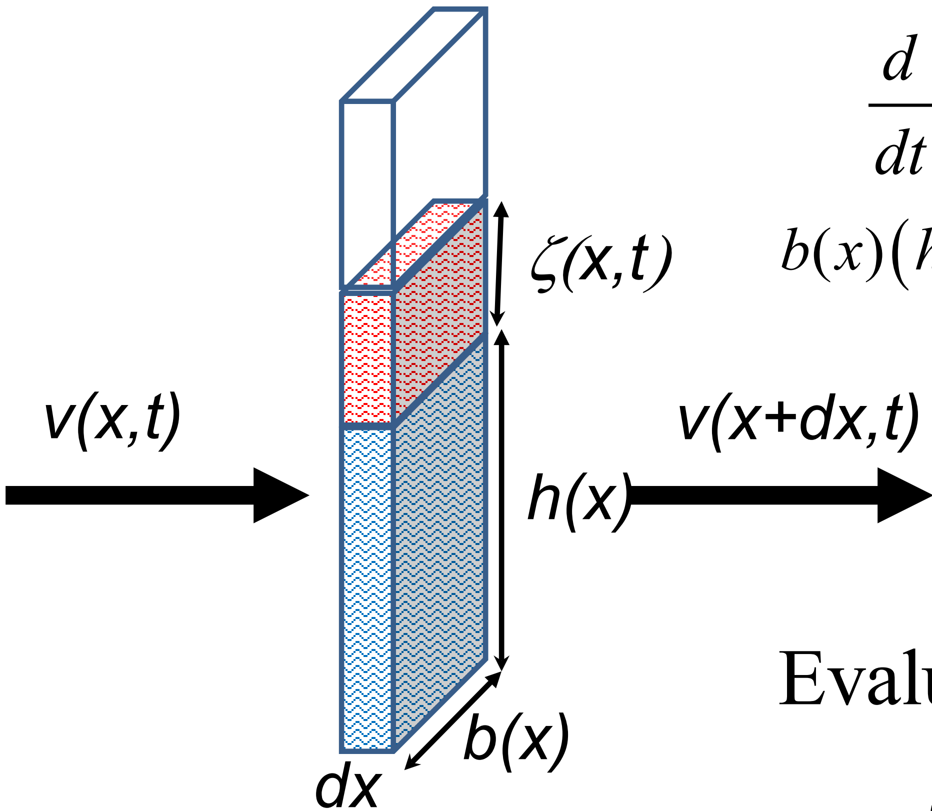
within the water

Horizontal fluid motions (keeping leading terms):

$$\frac{dv_x}{dt} \approx \frac{\partial v_x}{\partial t} = -\frac{1}{\rho}\frac{\partial p}{\partial x} = -g\frac{\partial \zeta}{\partial x}$$

$$\frac{dv_y}{dt} \approx \frac{\partial v_y}{\partial t} = -\frac{1}{\rho}\frac{\partial p}{\partial y} = -g\frac{\partial \zeta}{\partial y}$$

Consider a surface wave moving in the *x*-direction in a channel of width *b(x)* and height *h(x)* +*ζ(x,t)* :



Continuity condition in integral form:

$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = 0$$

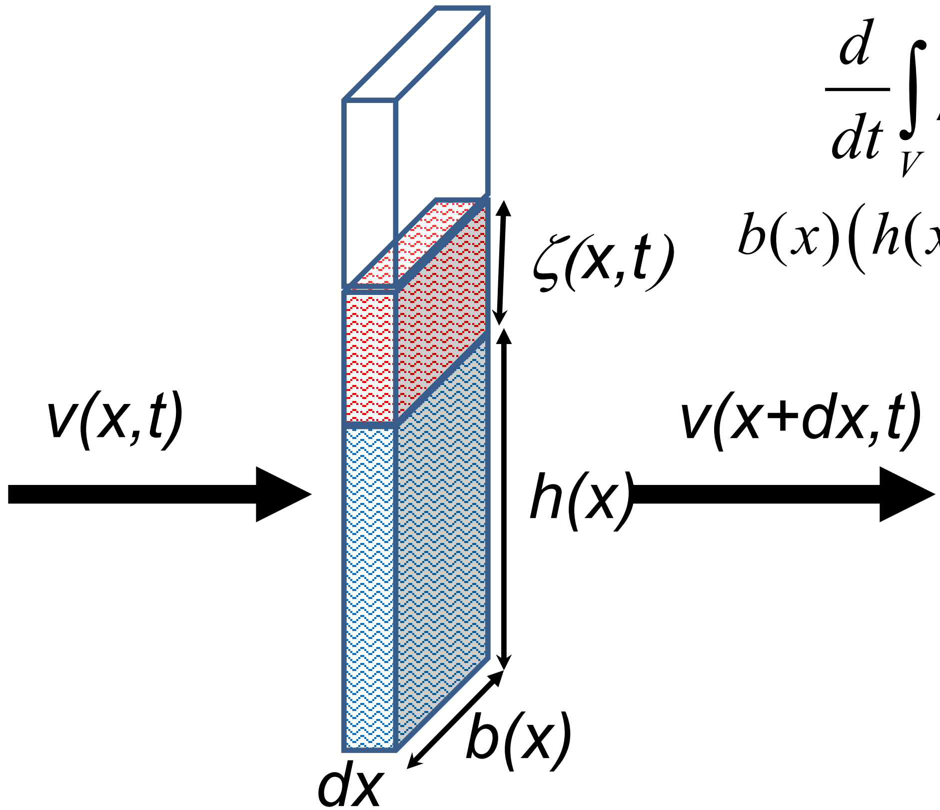
$$b(x)\big(h(x)+\zeta(x,t)\big)\hat{\mathbf{x}}$$

$$b(x)\big(h(x)+\zeta(x,t)\big)dx$$

Evaluating continuity condition:

$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\big(h(x)b(x)v(x,t)\big)$$

Some details



Continuity condition in integral form:

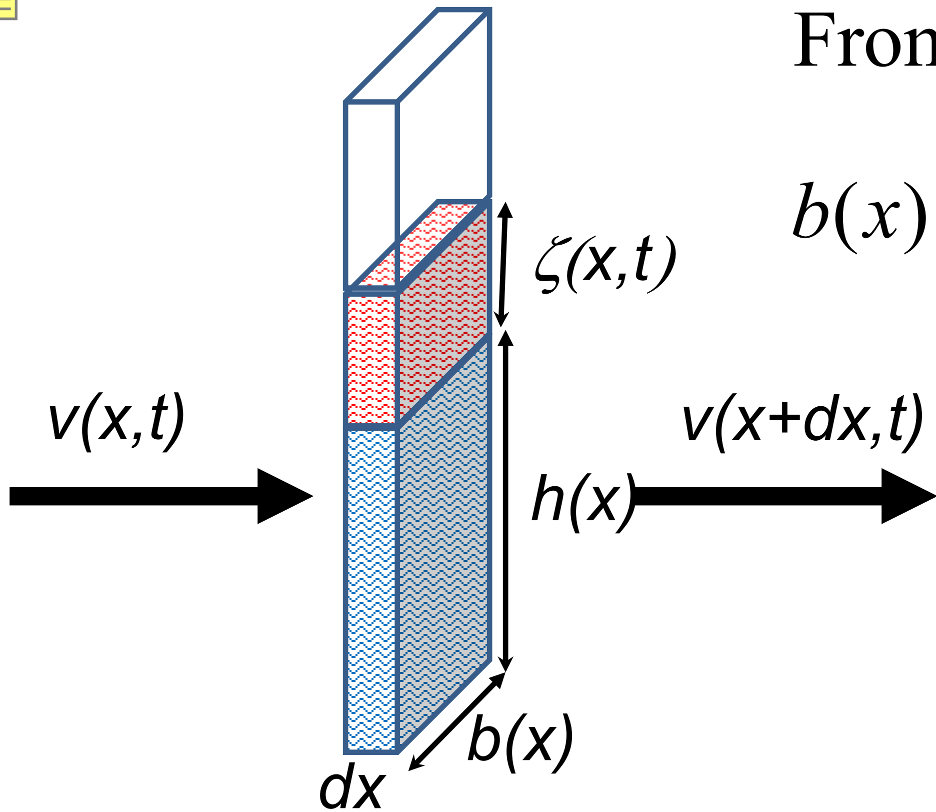
$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = 0$$

$$b(x)\left(h(x)+\zeta(x,t)\right)dx \qquad b(x)\left(h(x)+\zeta(x,t)\right)\hat{\mathbf{x}}$$

Here, we are assuming that ρ is constant

$$\frac{d}{dt}\int_V \rho dV + \int_A \rho \mathbf{v}\cdot d\mathbf{A} = \rho\int b(x)\frac{\partial \zeta}{\partial t}dx + \rho\int \frac{\partial}{\partial x}\left(b(x)(h(x)+\zeta(x,t))v(x,t)\right)dx = 0$$

$$\Rightarrow b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\left(h(x)b(x)v(x,t)\right)$$



From continuity condition:

$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\big(h(x)b(x)v(x,t)\big)$$

Example (Problem 10.3):

$$b(x) = b_0 \qquad h(x) = \kappa x$$

(A special case sometimes found at a beach.)

$$b_0 \frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\big((\kappa x) b_0 v(x,t)\big)$$

$$\frac{\partial \zeta}{\partial t} = -\kappa\left(v + x\frac{\partial v}{\partial x}\right)$$

From Newton-Euler equation:

$$\frac{dv}{dt} \approx \frac{\partial v}{\partial t} = -g\frac{\partial \zeta}{\partial x}$$

Example continued

$$\frac{\partial \zeta}{\partial t} = -\kappa\left(v + x\frac{\partial v}{\partial x}\right) \quad \Rightarrow \quad \frac{\partial^2 \zeta}{\partial t^2} = -\kappa\left(\frac{\partial v}{\partial t} + x\frac{\partial^2 v}{\partial x \partial t}\right)$$

$$\frac{\partial v}{\partial t} = -g\frac{\partial \zeta}{\partial x} \quad \Rightarrow \quad \frac{\partial^2 \zeta}{\partial t^2} = \kappa g\left(\frac{\partial \zeta}{\partial x} + x\frac{\partial^2 \zeta}{\partial x^2}\right)$$

It can be shown that a solution can take the form:

$$\zeta(x,t) = CJ_0\left(\frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}\right)\cos(\omega t)$$

Note that $J_0(u)$ satisfies the equation: $\left(\frac{d^2}{du^2} + \frac{1}{u}\frac{d}{du} + 1\right)J_0(u) = 0$

Therefore, for $u = \frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}$

$$\left(x\frac{d^2}{dx^2} + \frac{d}{dx}\right)J_0(u) = \frac{\omega^2}{\kappa g}\left(\frac{d^2}{du^2} + \frac{1}{u}\frac{d}{du}\right)J_0(u) = -\frac{\omega^2}{\kappa g}J_0(u)$$

Therefore, for $u=\frac{2\omega}{\sqrt{\kappa g}}\sqrt{x} \quad \Rightarrow \frac{1}{\sqrt{x}}=\frac{2\omega}{\sqrt{\kappa g}}\frac{1}{u}$

$$\left(x\frac{d^2}{dx^2}+\frac{d}{dx}\right)J_0(u)=\frac{\omega^2}{\kappa g}\left(\frac{d^2}{du^2}+\frac{1}{u}\frac{d}{du}\right)J_0(u)=-\frac{\omega^2}{\kappa g}J_0(u)$$

Detail: $\frac{dJ_0(u)}{dx}=\frac{dJ_0(u)}{du}\frac{\omega}{\sqrt{\kappa g}}\frac{1}{\sqrt{x}}$

$$\frac{d^2J_0(u)}{dx^2}=\frac{d^2J_0(u)}{du^2}\left(\frac{\omega}{\sqrt{\kappa g}}\frac{1}{\sqrt{x}}\right)^2-\frac{dJ_0(u)}{du}\frac{\omega}{2\sqrt{\kappa g}}\frac{1}{x\sqrt{x}}$$

Therefore: $\left(x\frac{d^2}{dx^2}+\frac{d}{dx}\right)J_0(u)=\left(\frac{\omega^2}{\kappa g}\frac{d^2J_0(u)}{du^2}+\frac{dJ_0(u)}{du}\frac{\omega}{2\sqrt{\kappa g}}\frac{1}{\sqrt{x}}\right)$

$$=\frac{\omega^2}{\kappa g}\left(\frac{d^2J_0(u)}{du^2}+\frac{dJ_0(u)}{du}\frac{1}{u}\right)$$

Example continued

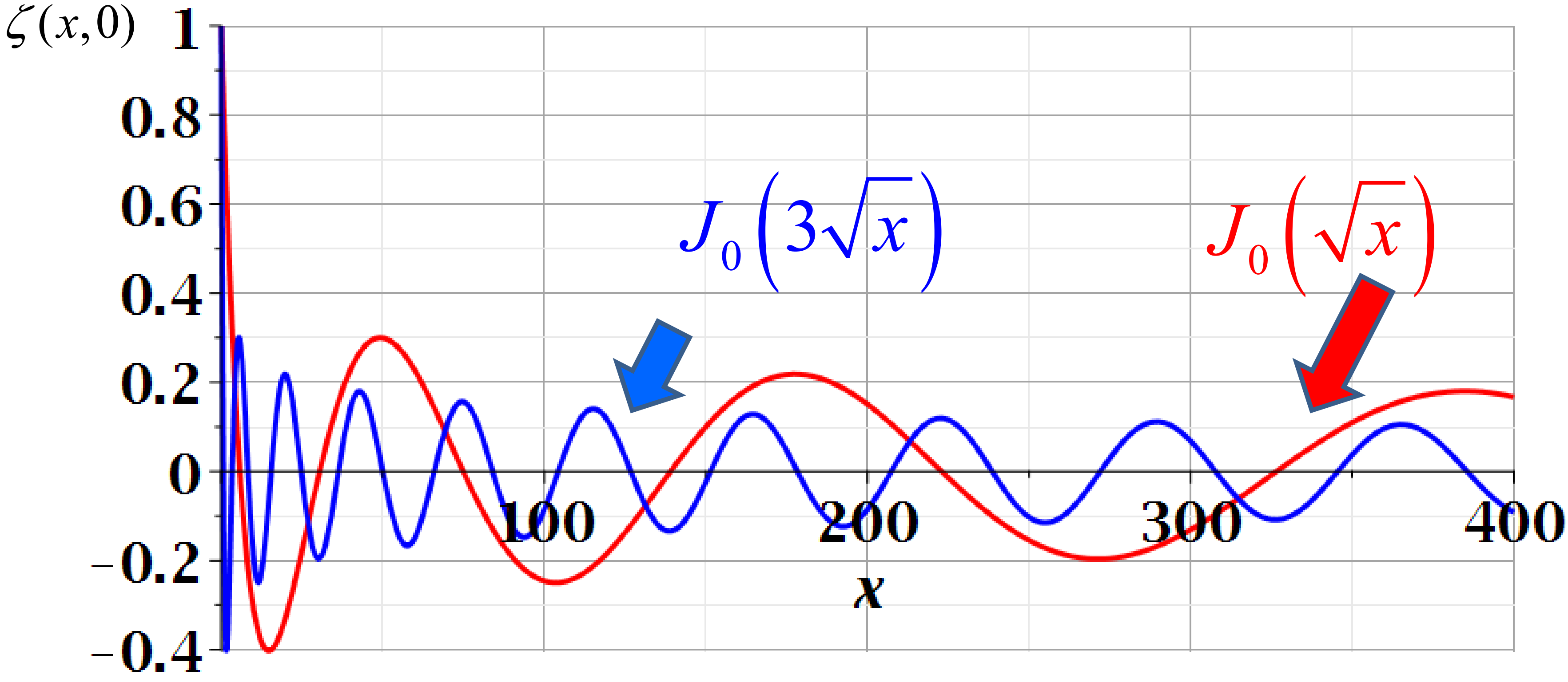
$$\frac{\partial^2 \zeta}{\partial t^2} = \kappa g \left(\frac{\partial \zeta}{\partial x} + x \frac{\partial^2 \zeta}{\partial x^2} \right)$$

$$\Rightarrow \zeta(x,t) = C J_0\left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t)$$

Check:

$$-\omega^2 C J_0\left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t) = \kappa g \left(\frac{\partial}{\partial x} + x \frac{\partial^2}{\partial x^2} \right) C J_0\left(\frac{2\omega\sqrt{x}}{\sqrt{\kappa g}} \right) \cos(\omega t)$$

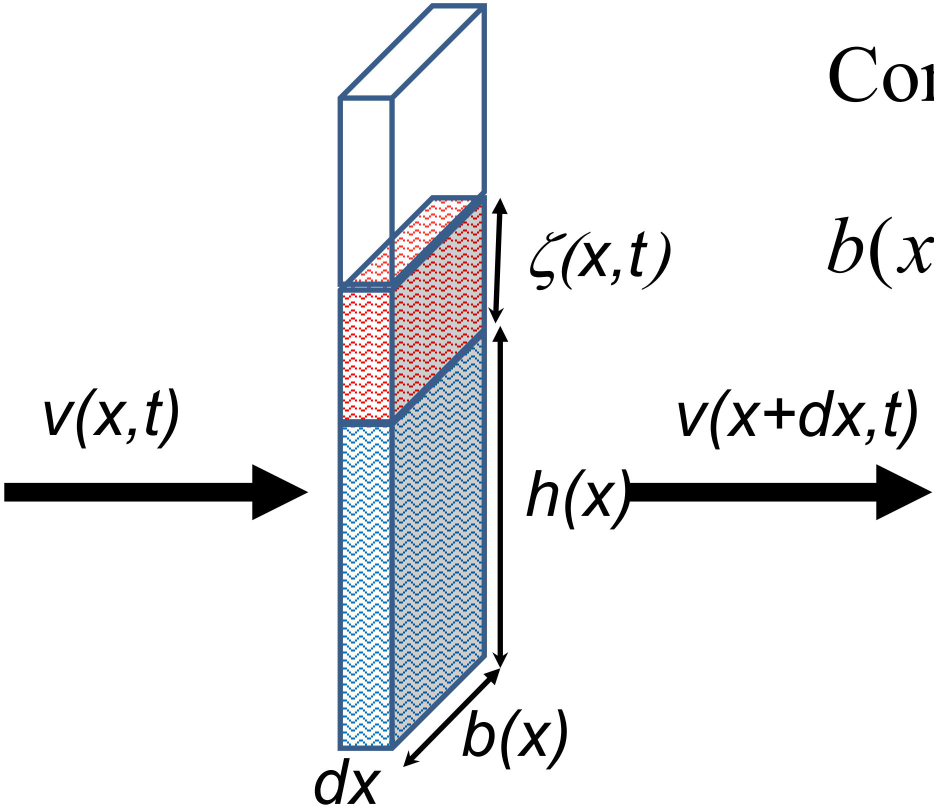
$$\zeta(x,t) = CJ_0\left(\frac{2\omega}{\sqrt{\kappa g}}\sqrt{x}\right)\cos(\omega t)$$



Imagine watching the waves at a beach – can you visualize the configuration for the surface wave pattern to approximation this situation?

a. Long flat beach
b. Beach in which average water level increases
c. Beach in which average water level decreases

A simplier example:



Continuity condition:

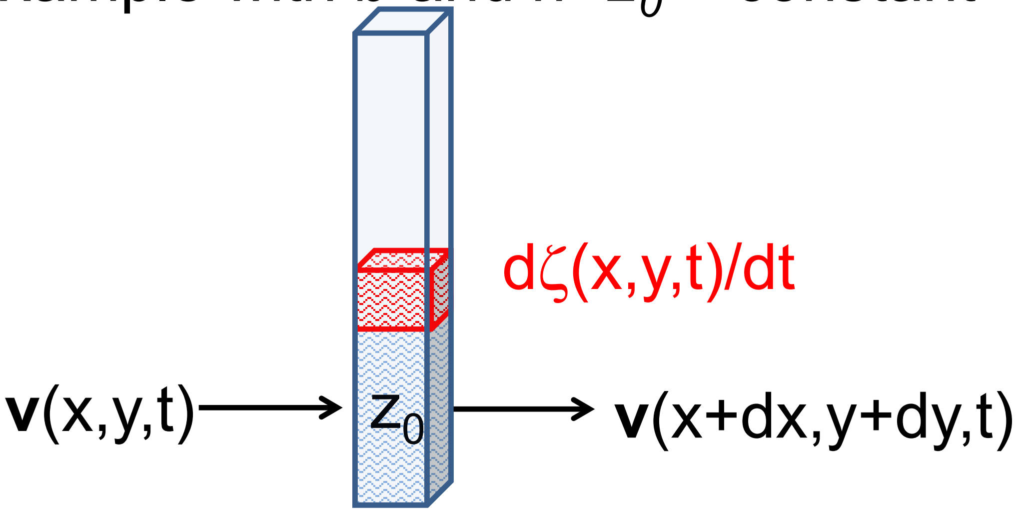
$$b(x)\frac{\partial \zeta}{\partial t} = -\frac{\partial}{\partial x}\big(h(x)b(x)v(x,t)\big)$$

Special case, where *b* and *h* are constant --

For constant b and h:

$$\frac{\partial \zeta}{\partial t} = -h\frac{\partial}{\partial x}\big(v(x,t)\big)$$

Example with b and $h=z_0=$ constant -- continued



Continuity condition for flow of incompressible fluid:

$$\frac{\partial \zeta}{\partial t} + h\nabla \cdot \mathbf{v} = 0$$

From horizontal flow relations: $\frac{\partial \mathbf{v}}{\partial t} = -g\nabla \zeta$

Equation for surface function: $\frac{\partial^2 \zeta}{\partial t^2} - gh\nabla^2 \zeta = 0$

For uniform channel:

Surface wave equation:

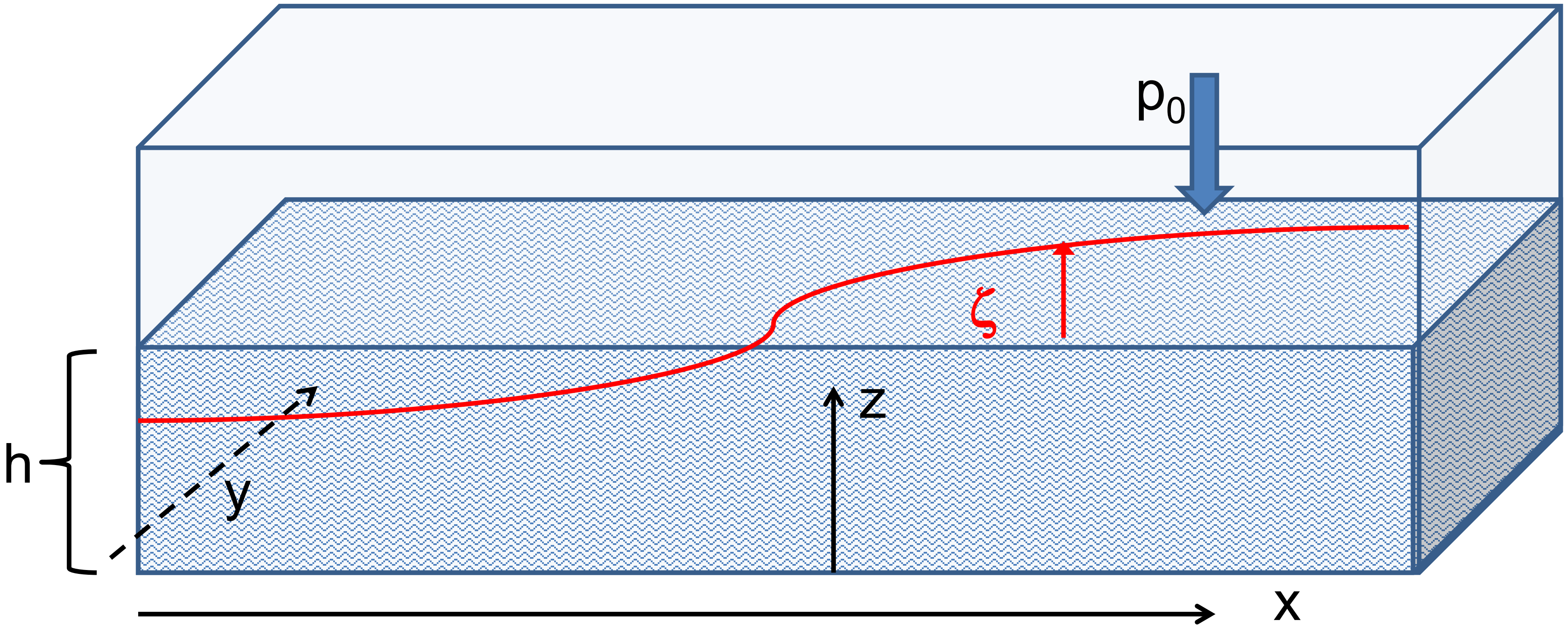
$$\frac{\partial^2 \zeta}{\partial t^2} - c^2 \nabla^2 \zeta = 0 \qquad c^2 = gh$$

More complete analysis finds:

$$c^2 = \frac{g}{k} \tanh(kh) \qquad \text{where } k = \frac{2\pi}{\lambda}$$

More details: -- recall setup --

Consider a container of water with average height h and surface h+$\zeta(x,y,t)$



Equations describing fluid itself (without boundaries)

Euler's equation for incompressible fluid:

$$\frac{d\mathbf{v}}{dt} = \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v}\cdot\nabla\mathbf{v} = \frac{\partial \mathbf{v}}{\partial t} + \nabla\left(\tfrac{1}{2}v^2\right) + \mathbf{v}\times\left(\nabla\times\mathbf{v}\right) = -\nabla U - \frac{\nabla p}{\rho}$$

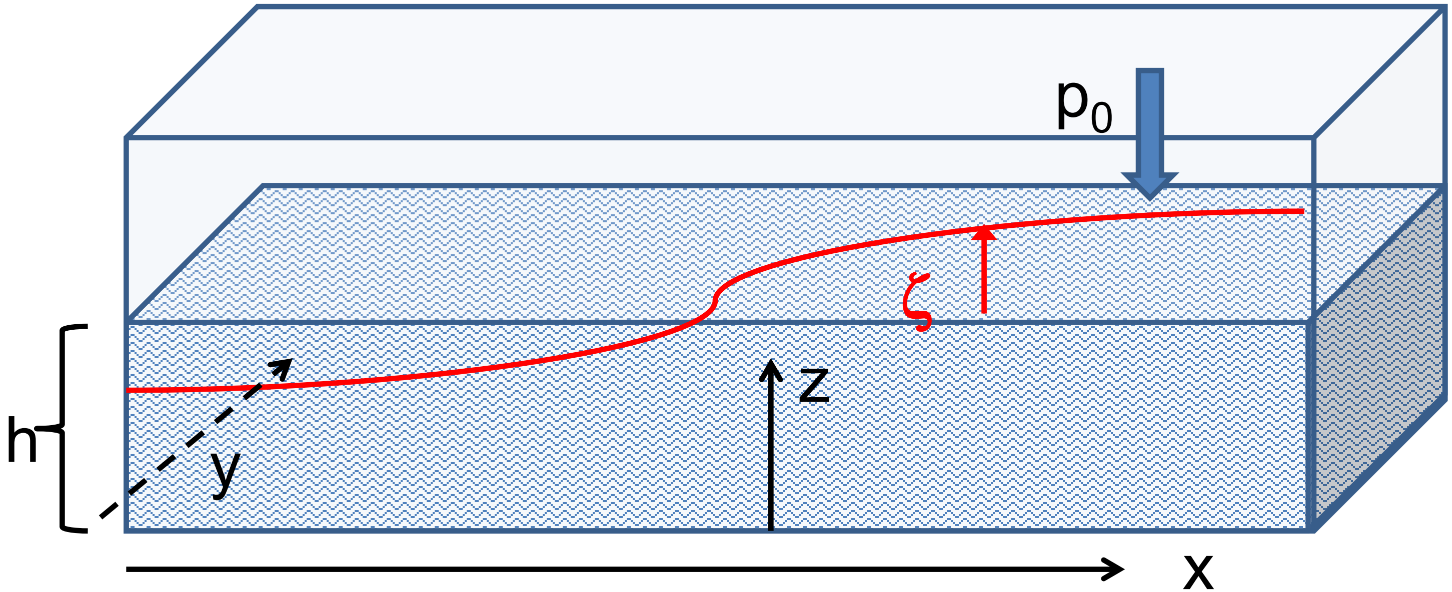
Assume that $\nabla\times\mathbf{v} = 0$ (irrotational flow) $\Rightarrow \mathbf{v} = -\nabla\Phi$

$$\Rightarrow \nabla\left(-\frac{\partial\Phi}{\partial t} + \tfrac{1}{2}v^2 + U + \frac{p}{\rho}\right) = 0$$

$$\Rightarrow -\frac{\partial\Phi}{\partial t} + \tfrac{1}{2}v^2 + U + \frac{p}{\rho} = \text{constant (within the fluid)}$$

For the same system, the continuity condition becomes

$$\nabla\cdot\mathbf{v} = -\nabla^2\Phi = 0$$



Within fluid: $0 \leq z \leq h + \zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2} v^2 + g\left(z - h\right) = \text{constant}$$

(We have absorbed p_0 in "constant")

$$-\nabla^2 \Phi = 0$$

At surface: $z = h + \zeta$ with $\zeta = \zeta\left(x, y, t\right)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y}$$

where $v_{x,y} = v_{x,y}\left(x, y, h + \zeta, t\right)$

Full equations:

Within fluid: $0 \le z \le h+\zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2}v^2 + g\left(z-h\right) = \text{constant}$$

(We have absorbed p_0 in "constant")

$$-\nabla^2 \Phi = 0$$

At surface: $z = h+\zeta$ with $\zeta = \zeta\left(x,y,t\right)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y} \qquad \text{where } v_{x,y} = v_{x,y}\left(x,y,h+\zeta,t\right)$$

Linearized equations:

$$\text{For} \quad 0 \le z \le h+\zeta: \quad -\frac{\partial \Phi}{\partial t} + g\left(z-h\right) = 0 \qquad -\nabla^2 \Phi = 0$$

$$\text{At surface}: \quad z = h+\zeta \qquad \frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} = v_z\left(x,y,h+\zeta,t\right)$$

$$-\frac{\partial \Phi\left(x,y,h+\zeta,t\right)}{\partial t} + g\zeta = 0$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x*:

For $\quad 0 \le z \le h + \zeta : \qquad \nabla^2 \Phi = \left(\frac{\partial^2}{\partial z^2} + \frac{\partial^2}{\partial x^2} \right) \Phi(x,z,t) = 0$

Consider and periodic waveform: $\quad \Phi(x,z,t) = Z(z)\cos\big(k(x-ct)\big)$

$$\Rightarrow \left(\frac{d^2}{dz^2} - k^2 \right) Z(z) = 0$$

Boundary condition at bottom of tank: $\quad v_z(x,0,t) = 0$

$$\Rightarrow \frac{dZ}{dz}(0) = 0 \qquad\qquad Z(z) = A\cosh(kz)$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$\text{At surface:} \quad z = h+\zeta \qquad \frac{\partial \zeta}{\partial t} = v_z\left(x, h+\zeta, t\right) = -\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial z}$$

$$-\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial t} + g\zeta = 0$$

$$-\frac{\partial^2 \Phi\left(x, h+\zeta, t\right)}{\partial t^2} + g\frac{\partial \zeta}{\partial t} = -\frac{\partial^2 \Phi\left(x, h+\zeta, t\right)}{\partial t^2} - g\frac{\partial \Phi\left(x, h+\zeta, t\right)}{\partial z} = 0$$

$$\text{For} \qquad \Phi(x,\left(h+\zeta\right),t) = A\cosh(k\left(h+\zeta\right))\cos\left(k(x-ct)\right)$$

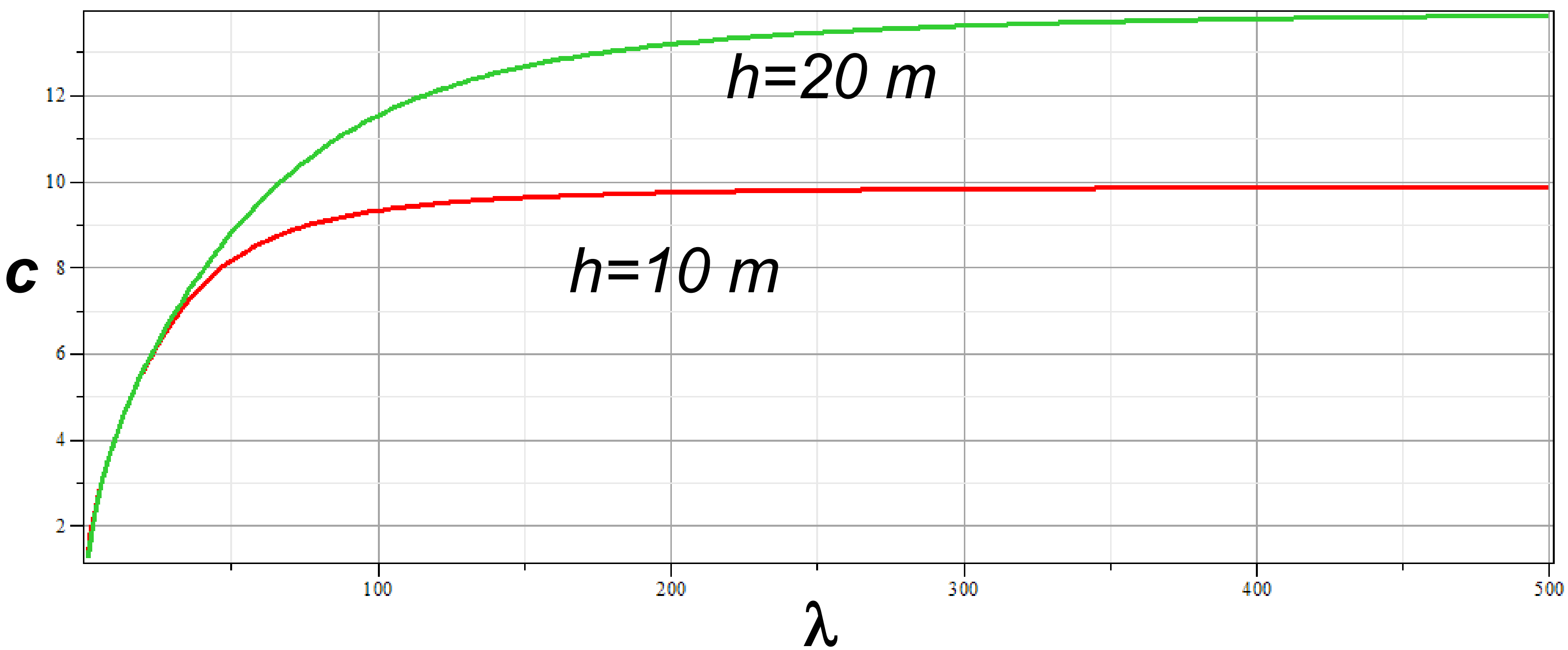
$$A\cosh(k\left(h+\zeta\right))\cos\left(k(x-ct)\right)\left(k^2c^2 - gk\frac{\sinh(k\left(h+\zeta\right))}{\cosh(k\left(h+\zeta\right))}\right) = 0$$

$$\Rightarrow \mathrm{c}^2 = \frac{g}{k}\frac{\sinh(k\left(h+\zeta\right))}{\cosh(k\left(h+\zeta\right))}$$

For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$c^2 = \frac{g}{k}\frac{\sinh(k(h+\zeta))}{\cosh(k(h+\zeta))} = \frac{g}{k}\tanh(k(h+\zeta))$$

Assuming $\zeta << h$: $\quad c^2 = \frac{g}{k}\tanh(kh) \quad\quad \lambda = \frac{2\pi}{k}$



For simplicity, keep only linear terms and assume that horizontal variation is only along *x* – continued:

$$c^2 \approx \frac{g}{k}\tanh(kh) \qquad \text{For } \lambda >> h,\ c^2 \approx gh$$

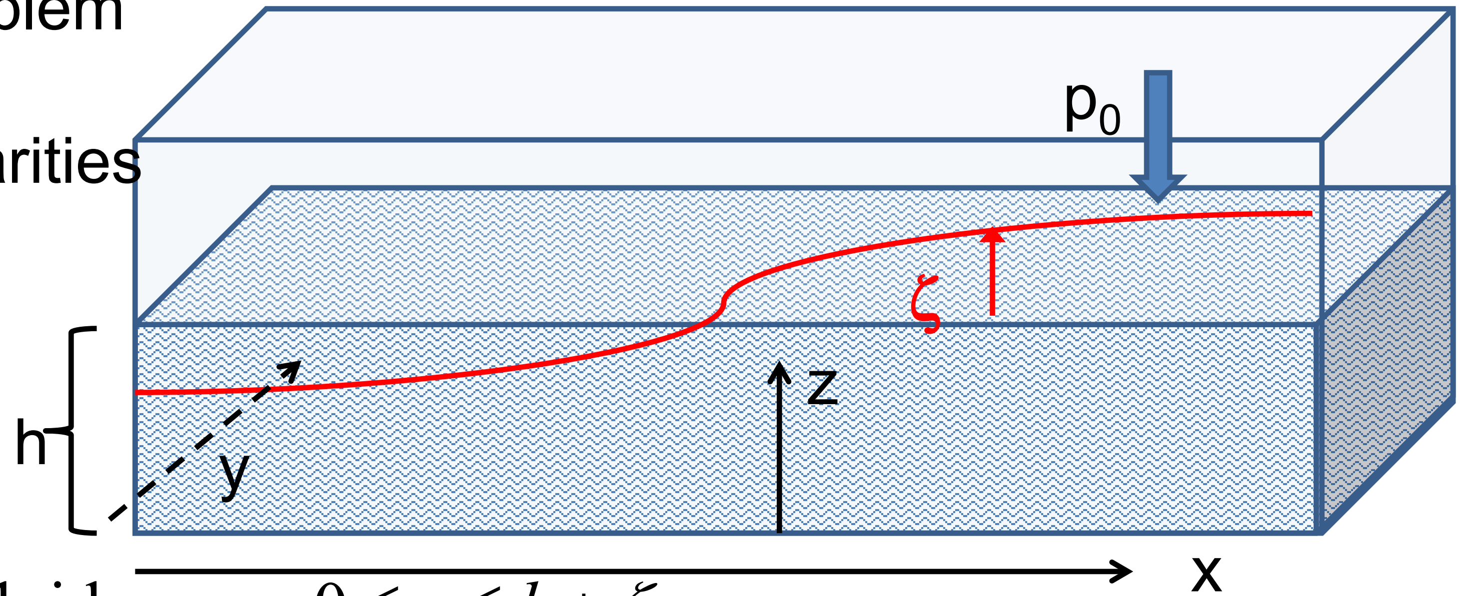
$$\Phi(x,z,t) = A\cosh(kz)\cos(k(x-ct))$$

$$\zeta(x,t) = \frac{1}{g}\frac{\partial\Phi(x,h+\zeta,t)}{\partial t} \approx \frac{kc}{g}A\cosh(kh)\sin(k(x-ct))$$

Note that for $\lambda >> h, \quad c^2 \approx gh$

(solutions are consistent with previous analysis)

General problem including non-linearities



Within fluid: $\quad 0 \le z \le h + \zeta$

$$-\frac{\partial \Phi}{\partial t} + \tfrac{1}{2} v^2 + g(z-h) = \text{constant} \quad \text{(We have absorbed}$$

$$-\nabla^2 \Phi = 0 \qquad\qquad p_0 \text{ in our constant.)}$$

At surface: $\quad z = h + \zeta \qquad$ with $\zeta = \zeta(x, y, t)$

$$\frac{d\zeta}{dt} = \frac{\partial \zeta}{\partial t} + v_x \frac{\partial \zeta}{\partial x} + v_y \frac{\partial \zeta}{\partial y} \qquad \text{where } v_{x,y} = v_{x,y}(x, y, h+\zeta, t)$$