



PHY 711 Classical Mechanics and Mathematical Methods

10-10:50 AM MWF in Olin 103

Notes on Lecture 5 – Chap. 3&6(33) in F&W

Lagrangian mechanics

1. Lagrange's equations in the presence of velocity dependent potentials – such as obtained when a charged particle moves in a magnetic field.
2. Effects of constraints

Your questions –

From Thomas –

On slide 28 I am not sure how you combined the coupled equations? On slide 31 I think there were some steps that were not shown. I think I get what happened, but I would need to work it out.

From Julia --

Would there ever be a situation where it would make sense to find the Lorentz force in non-Cartesian coordinates?

From Conall --

I am not understanding how we come to the conclusion on slide 11. I was also looking through 6.33 in the book and am equally confused with that.

Physics Colloquium

- Thursday -
September 5,
2024

Welcome

Professor Daniel Kim-Shapiro
Professor Paul Anderson

Summer Research Presentations

Department Introductions and
Announcements

Welcome!

Department of
Physics



WAKE FOREST
UNIVERSITY

4 PM Olin 101



Course schedule

(Preliminary schedule -- subject to frequent adjustment.)

| | Date | F&W | Topic | HW |
|---|----------------|-------------|-------------------------------------|--------------------|
| 1 | Mon, 8/26/2024 | | Introduction and overview | #1 |
| 2 | Wed, 8/28/2024 | Chap. 3(17) | Calculus of variation | #2 |
| 3 | Fri, 8/30/2024 | Chap. 3(17) | Calculus of variation | #3 |
| 4 | Mon, 9/02/2024 | Chap. 3 | Lagrangian equations of motion | #4 |
| 5 | Wed, 9/04/2024 | Chap. 3 & 6 | Lagrangian equations of motion | #5 |
| 6 | Fri, 9/06/2024 | Chap. 3 & 6 | Lagrangian equations of motion | |
| 7 | Mon, 9/09/2024 | Chap. 3 & 6 | Lagrangian to Hamiltonian formalism | |



PHY 711 – Assignment #5

Assigned: 09/04/2024 Due: 09/09/2024

A particle of mass m and charge q is subjected to a vector potential $\mathbf{A}(\mathbf{r}, t) = -(E_0ct + B_0x)\hat{\mathbf{z}}$. In this case, the scalar potential is zero: $\Phi(\mathbf{r}, t) = 0$. (Note that we are using the cgs Gaussian units of your text book.) Here E_0 denotes a constant electric field amplitude and B_0 denotes a constant magnetic field amplitude. The initial particle position is $\mathbf{r}(t = 0) = 0$ and the initial particle velocity is $\dot{\mathbf{r}}(t = 0) = 0$.

- a. Determine the Lagrangian $L(x, y, z, \dot{x}, \dot{y}, \dot{z}, t)$ which describes the particle's motion.
- b. Write the Euler-Lagrange equations for this system.
- c. Find and evaluate the constants of motion for this system.
- d. Find the particle trajectories $x(t)$, $y(t)$, $z(t)$ by solving the equations and imposing the given initial conditions.

Comment on single and multiple coordinates --

Hamilton's principle for optimization for a single trajectory $q(t)$:

$$S = \int_{t_i}^{t_f} L(q, \dot{q}, t) dt \quad \text{where } L(q, \dot{q}, t) = \text{Kinetic energy} - \text{Potential energy}$$

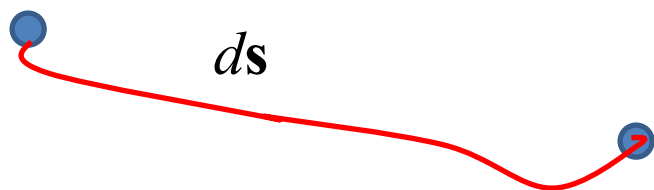
Hamilton's principle for optimization for a multiple trajectories $\{q_\sigma(t)\}$:

$$S = \int_{t_i}^{t_f} L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) dt \quad \text{where } L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) = \text{Kinetic energy} - \text{Potential energy}$$

This "works" provided that the variation of each trajectory $q_\sigma(t)$ can be analyzed.

Note that the trajectory components can be independent (as in the case of cartesian coordinates and/or multiple particles or can be dependent in which case we can use the "trick" of Lagrange multipliers.

Previously derived form for the Lagrangian --



Generalized coordinates :

$$q_{\sigma}(\{x_i\})$$

$$(\mathbf{F} - m\mathbf{a}) \cdot d\mathbf{s} = -\sum_{\sigma} \frac{\partial U}{\partial q_{\sigma}} \delta q_{\sigma} - \sum_{\sigma} \left(\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_{\sigma}} - \frac{\partial T}{\partial q_{\sigma}} \right) \delta q_{\sigma} = 0$$

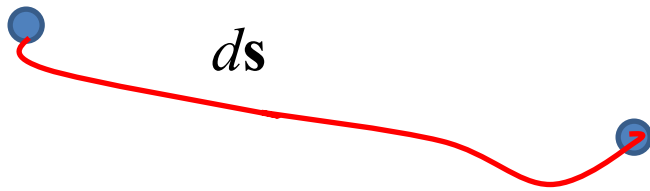
$$= -\sum_{\sigma} \left(\frac{d}{dt} \frac{\partial (T - U)}{\partial \dot{q}_{\sigma}} - \frac{\partial (T - U)}{\partial q_{\sigma}} \right) \delta q_{\sigma} = 0$$

$$= -\sum_{\sigma} \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_{\sigma}} - \frac{\partial L}{\partial q_{\sigma}} \right) \delta q_{\sigma} = 0$$

$$L(\{q_{\sigma}\}, \{\dot{q}_{\sigma}\}; t) = T - U$$

Note: This is only true if

$$\frac{\partial U}{\partial \dot{q}_{\sigma}} = 0$$



Generalized coordinates :
 $q_\sigma(\{x_i\})$

Define -- Lagrangian: $L \equiv T - U$

$$L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t)$$

$$(\mathbf{F}-m\mathbf{a}) \cdot d\mathbf{s} = - \sum_\sigma \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} \right) \delta q_\sigma = 0$$

$$\Rightarrow \text{Minimization integral: } S = \int_{t_i}^{t_f} L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) dt$$

→ Hamilton's principle from the "backwards" application of the Euler-Lagrange equations to

Define -- Lagrangian: $L \equiv T - U$

$$L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t)$$



Summary –

Hamilton's principle:

Given the Lagrangian function: $L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) \equiv T - U,$

The physical trajectories of the generalized coordinates $\{q_\sigma(t)\}$

are those which minimize the action: $S = \int L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) dt$

Euler-Lagrange equations:

$$\sum_{\sigma} \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_{\sigma}} - \frac{\partial L}{\partial q_{\sigma}} \right) \delta q_{\sigma} = 0 \quad \Rightarrow \text{for each } \sigma : \quad \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_{\sigma}} - \frac{\partial L}{\partial q_{\sigma}} \right) = 0$$

Note: in “proof” of Hamilton’s principle:

$$\left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} \right) = 0 \quad \text{for} \quad L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) \equiv T - U$$

It was necessary to assume that :

$\frac{d}{dt} \frac{\partial U}{\partial \dot{q}_\sigma}$ does not contribute to the result.

⇒ How can we represent velocity-dependent forces?

Why do we need velocity dependent forces?

- a. Friction is sometimes represented as a velocity dependent force. (difficult to treat with Lagrangian mechanics.)
- b. Lorentz force on a moving charged particle in the presence of a magnetic field.

Some details --

$$\left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} \right) = 0 \quad \text{for} \quad L = L(\{q_\sigma\}, \{\dot{q}_\sigma\}, t) \equiv T - U$$

It was necessary to assume that:

$$\frac{d}{dt} \frac{\partial U}{\partial \dot{q}_\sigma} \quad \text{does not contribute to the result.}$$

This comes from D'Alembert's analysis which gave us:

$$(\mathbf{F} - m\mathbf{a}) \cdot d\mathbf{s} = 0 = - \sum_{\sigma} \frac{\partial U}{\partial q_{\sigma}} \delta q_{\sigma} - \sum_{\sigma} \left(\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_{\sigma}} - \frac{\partial T}{\partial q_{\sigma}} \right) \delta q_{\sigma}$$

$$(\mathbf{F} - m\mathbf{a}) \cdot d\mathbf{s} = 0 = - \sum_{\sigma} \left(\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_{\sigma}} - \frac{\partial(T - U)}{\partial q_{\sigma}} \right) \delta q_{\sigma}$$

$$\text{while we want to use: } 0 = - \sum_{\sigma} \left(\frac{d}{dt} \frac{\partial(T - U)}{\partial \dot{q}_{\sigma}} - \frac{\partial(T - U)}{\partial q_{\sigma}} \right) \delta q_{\sigma}$$

Lorentz forces:

For particle of charge q in an electric field $\mathbf{E}(\mathbf{r}, t)$ and magnetic field $\mathbf{B}(\mathbf{r}, t)$:

$$\text{Lorentz force: } \mathbf{F} = q\left(\mathbf{E} + \frac{1}{c} \mathbf{v} \times \mathbf{B}\right)$$

x -component: $F_x = q\left(E_x + \frac{1}{c}(\mathbf{v} \times \mathbf{B})_x\right)$
In this case, it is convenient to use cartesian coordinates

$$L = L(x, y, z, \dot{x}, \dot{y}, \dot{z}, t) \equiv T - U$$

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

$$x\text{-component: } \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} \right) = 0$$

$$\text{Apparently: } F_x = -\frac{\partial U}{\partial x} + \frac{d}{dt} \frac{\partial U}{\partial \dot{x}}$$

$$\text{Answer: } U = q\Phi(\mathbf{r}, t) - \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

$$\text{where } \mathbf{E}(\mathbf{r}, t) = -\nabla\Phi(\mathbf{r}, t) - \frac{1}{c} \frac{\partial \mathbf{A}(\mathbf{r}, t)}{\partial t}$$

$$\mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$$

Note: Here we are using cartesian coordinates for convenience.

Note: We need to prove this!

More details --

$$\text{Consider: } 0 = -\sum_{\sigma} \left(\frac{d}{dt} \frac{\partial(T-U)}{\partial \dot{q}_{\sigma}} - \frac{\partial(T-U)}{\partial q_{\sigma}} \right) \delta q_{\sigma}$$

$$\text{Suppose } T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

$$\Rightarrow 0 = \left(\frac{d}{dt} \frac{\partial(T-U)}{\partial \dot{x}} - \frac{\partial(T-U)}{\partial x} \right) = \frac{d}{dt} m \dot{x} - \frac{d}{dt} \left(\frac{\partial U}{\partial \dot{x}} \right) + \frac{\partial U}{\partial x}$$

$$\Rightarrow m \ddot{x} = \frac{d}{dt} \left(\frac{\partial U}{\partial \dot{x}} \right) - \frac{\partial U}{\partial x} = F_x$$



Can be used to represent Lorentz force.

Units for electromagnetic fields and forces

cgs Gaussian units -- (as used your textbook)

E and **B** fields as related to vector and scalar potentials:

$$\mathbf{E}(\mathbf{r}, t) = -\nabla\Phi(\mathbf{r}, t) - \frac{1}{c} \frac{\partial \mathbf{A}(\mathbf{r}, t)}{\partial t} \quad \mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$$

Corresponding Lagrangian potential:

$$U = q\Phi(\mathbf{r}, t) - \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

SI units --

E and **B** fields as related to vector and scalar potentials:

$$\mathbf{E}(\mathbf{r}, t) = -\nabla\Phi(\mathbf{r}, t) - \frac{\partial \mathbf{A}(\mathbf{r}, t)}{\partial t} \quad \mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$$

Corresponding Lagrangian potential:

$$U = q\Phi(\mathbf{r}, t) - q\dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

Lorentz forces, continued:

x – component of Lorentz force: $F_x = q(E_x + \frac{1}{c}(\mathbf{v} \times \mathbf{B})_x)$

Suppose:
$$U = q\Phi(\mathbf{r}, t) - \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

Consider:
$$F_x = -\frac{\partial U}{\partial x} + \frac{d}{dt} \frac{\partial U}{\partial \dot{x}}$$

$$-\frac{\partial U}{\partial x} = -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} + \frac{q}{c} \left(\dot{x} \frac{\partial A_x(\mathbf{r}, t)}{\partial x} + \dot{y} \frac{\partial A_y(\mathbf{r}, t)}{\partial x} + \dot{z} \frac{\partial A_z(\mathbf{r}, t)}{\partial x} \right)$$

$$\frac{\partial U}{\partial \dot{x}} = -\frac{q}{c} A_x(\mathbf{r}, t)$$

$$\frac{d}{dt} \frac{\partial U}{\partial \dot{x}} = -\frac{q}{c} \frac{dA_x(\mathbf{r}, t)}{dt} = -\frac{q}{c} \left(\frac{\partial A_x(\mathbf{r}, t)}{\partial x} \dot{x} + \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \dot{y} + \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \dot{z} + \frac{\partial A_x(\mathbf{r}, t)}{\partial t} \right)$$

Lorentz forces, continued:

$$-\frac{\partial U}{\partial x} = -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} + \frac{q}{c} \left(\dot{x} \frac{\partial A_x(\mathbf{r}, t)}{\partial x} + \dot{y} \frac{\partial A_y(\mathbf{r}, t)}{\partial x} + \dot{z} \frac{\partial A_z(\mathbf{r}, t)}{\partial x} \right)$$

$$\frac{d}{dt} \frac{\partial U}{\partial \dot{x}} = -\frac{q}{c} \left(\frac{\partial A_x(\mathbf{r}, t)}{\partial x} \dot{x} + \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \dot{y} + \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \dot{z} + \frac{\partial A_x(\mathbf{r}, t)}{\partial t} \right)$$

$$\begin{aligned} F_x &= -\frac{\partial U}{\partial x} + \frac{d}{dt} \frac{\partial U}{\partial \dot{x}} \\ &= -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} + \frac{q}{c} \dot{y} \left(\frac{\partial A_y(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \right) + \frac{q}{c} \dot{z} \left(\frac{\partial A_z(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \right) - \frac{q}{c} \frac{\partial A_x(\mathbf{r}, t)}{\partial t} \\ &= -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} - \frac{q}{c} \frac{\partial A_x(\mathbf{r}, t)}{\partial t} + \frac{q}{c} \dot{y} \left(\frac{\partial A_y(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \right) + \frac{q}{c} \dot{z} \left(\frac{\partial A_z(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \right) \\ &= qE_x(\mathbf{r}, t) + \frac{q}{c} (\dot{y}B_z(\mathbf{r}, t) - \dot{z}B_y(\mathbf{r}, t)) = qE_x(\mathbf{r}, t) + \frac{q}{c} (\mathbf{v} \times \mathbf{B}(\mathbf{r}, t))_x \end{aligned}$$

Some details on last step:

$$\begin{aligned} F_x &= -\frac{\partial U}{\partial x} + \frac{d}{dt} \frac{\partial U}{\partial \dot{x}} \\ &= -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} + \frac{q}{c} \dot{y} \left(\frac{\partial A_y(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \right) + \frac{q}{c} \dot{z} \left(\frac{\partial A_z(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \right) - \frac{q}{c} \frac{\partial A_x(\mathbf{r}, t)}{\partial t} \\ &= -q \frac{\partial \Phi(\mathbf{r}, t)}{\partial x} - \frac{q}{c} \frac{\partial A_x(\mathbf{r}, t)}{\partial t} + \frac{q}{c} \dot{y} \left(\frac{\partial A_y(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial y} \right) + \frac{q}{c} \dot{z} \left(\frac{\partial A_z(\mathbf{r}, t)}{\partial x} - \frac{\partial A_x(\mathbf{r}, t)}{\partial z} \right) \end{aligned}$$

Note that: $\mathbf{E}(\mathbf{r}, t) = -\nabla \Phi(\mathbf{r}, t) - \frac{1}{c} \frac{\partial \mathbf{A}(\mathbf{r}, t)}{\partial t}$ $\mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$

So that:

$$F_x(\mathbf{r}, t) = qE_x(\mathbf{r}, t) + \frac{q}{c} (\dot{y}B_z(\mathbf{r}, t) - \dot{z}B_y(\mathbf{r}, t)) = qE_x(\mathbf{r}, t) + \frac{q}{c} (\mathbf{v} \times \mathbf{B}(\mathbf{r}, t))_x$$

It follows that similar analyses can be applied to $F_y(\mathbf{r}, t)$ and $F_z(\mathbf{r}, t)$.

Lorentz forces, continued:

Summary of results (using cartesian coordinates)

$$L = L(x, y, z, \dot{x}, \dot{y}, \dot{z}, t) \equiv T - U$$

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) \quad U = q\Phi(\mathbf{r}, t) - \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

$$\text{where } \mathbf{E}(\mathbf{r}, t) = -\nabla\Phi(\mathbf{r}, t) - \frac{1}{c} \frac{\partial \mathbf{A}(\mathbf{r}, t)}{\partial t} \quad \mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - q\Phi(\mathbf{r}, t) + \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

Note that, more generally, $U = U_{\text{mechanical}} + U_{EM}$

$$U_{EM} = q\Phi(\mathbf{r}, t) - \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

Example Lorentz force

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - q\Phi(\mathbf{r}, t) + \frac{q}{c} \dot{\mathbf{r}} \cdot \mathbf{A}(\mathbf{r}, t)$$

$$\text{Suppose } \mathbf{E}(\mathbf{r}, t) \equiv 0, \quad \mathbf{B}(\mathbf{r}, t) \equiv B_0 \hat{\mathbf{z}}$$

$$\mathbf{A}(\mathbf{r}, t) = \frac{1}{2} B_0 (-y \hat{\mathbf{x}} + x \hat{\mathbf{y}})$$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{q}{2c} B_0 (-\dot{x}y + \dot{y}x)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0 \quad \Rightarrow \quad \frac{d}{dt} \left(m\dot{x} - \frac{q}{2c} B_0 y \right) - \frac{q}{2c} B_0 \dot{y} = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} = 0 \quad \Rightarrow \quad \frac{d}{dt} \left(m\dot{y} + \frac{q}{2c} B_0 x \right) + \frac{q}{2c} B_0 \dot{x} = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{z}} - \frac{\partial L}{\partial z} = 0 \quad \Rightarrow \quad \frac{d}{dt} m\dot{z} = 0$$

Example Lorentz force -- continued

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{q}{2c} B_0 (-\dot{x}y + \dot{y}x)$$

$$\frac{d}{dt} \left(m\dot{x} - \frac{q}{2c} B_0 y \right) - \frac{q}{2c} B_0 \dot{y} = 0 \quad \Rightarrow \quad m\ddot{x} - \frac{q}{c} B_0 \dot{y} = 0$$

$$\frac{d}{dt} \left(m\dot{y} + \frac{q}{2c} B_0 x \right) + \frac{q}{2c} B_0 \dot{x} = 0 \quad \Rightarrow \quad m\ddot{y} + \frac{q}{c} B_0 \dot{x} = 0$$

$$\frac{d}{dt} m\dot{z} = 0 \quad \Rightarrow \quad m\ddot{z} = 0$$



Example Lorentz force -- continued

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{q}{2c} B_0 (-\dot{x}y + \dot{y}x)$$

$$m\ddot{x} = +\frac{q}{c} B_0 \dot{y}$$

$$m\ddot{y} = -\frac{q}{c} B_0 \dot{x}$$

$$m\ddot{z} = 0$$

Note that same equations are obtained from direct application of Newton's laws :

$$m\ddot{\mathbf{r}} = \frac{q}{c} \dot{\mathbf{r}} \times B_0 \hat{\mathbf{z}}$$

Example Lorentz force -- continued

Evaluation of equations :

$$m\ddot{x} - \frac{q}{c} B_0 \dot{y} = 0 \quad \dot{x}(t) = V_0 \sin\left(\frac{qB_0}{mc}t + \phi\right)$$

$$m\ddot{y} + \frac{q}{c} B_0 \dot{x} = 0 \quad \dot{y}(t) = V_0 \cos\left(\frac{qB_0}{mc}t + \phi\right)$$

$$m\ddot{z} = 0 \quad \dot{z}(t) = V_{0z}$$

$$x(t) = x_0 - \frac{mc}{qB_0} V_0 \cos\left(\frac{qB_0}{mc}t + \phi\right)$$

$$y(t) = y_0 + \frac{mc}{qB_0} V_0 \sin\left(\frac{qB_0}{mc}t + \phi\right)$$

$$z(t) = z_0 + V_{0z}t$$

Note: 3 second order differential equations need 6 constants for specific solution.

Some details

Evaluation of equations :

$$m\ddot{x} - \frac{q}{c} B_0 \dot{y} = 0 \qquad \frac{d}{dt} \left(\dot{x} - \frac{qB_0}{mc} y \right) = 0 \quad \Rightarrow \quad \dot{x} - \frac{qB_0}{mc} y = K_1$$

$$m\ddot{y} + \frac{q}{c} B_0 \dot{x} = 0 \qquad \frac{d}{dt} \left(\dot{y} + \frac{qB_0}{mc} x \right) = 0 \quad \Rightarrow \quad \dot{y} + \frac{qB_0}{mc} x = K_2$$

$$m\ddot{z} = 0 \qquad \frac{d}{dt} \dot{z} = 0 \qquad \Rightarrow \quad \dot{z} = K_3$$

How can you solve coupled differential equations?
How can you determine the constants K_1 , K_2 , K_3 ?

Example Lorentz force -- continued

Consider formulation with different Gauge: $\mathbf{A}(\mathbf{r}) = -B_0 y \hat{\mathbf{x}}$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - \frac{q}{c} B_0 \dot{x} y$$

$$\frac{d}{dt} \left(m \dot{x} - \frac{q}{c} B_0 y \right) = 0 \quad \Rightarrow \quad m \ddot{x} - \frac{q}{c} B_0 \dot{y} = 0$$

$$\frac{d}{dt} (m \dot{y}) + \frac{q}{c} B_0 \dot{x} = 0 \quad \Rightarrow \quad m \ddot{y} + \frac{q}{c} B_0 \dot{x} = 0$$

$$\frac{d}{dt} m \dot{z} = 0 \quad \Rightarrow \quad m \ddot{z} = 0$$

Does it surprise you that the same equations of motion are obtained with a different Gauge?

How do these two different forms of \mathbf{A} correspond to the same \mathbf{B} ?

$$\mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$$

Consider $\mathbf{A}'(\mathbf{r}, t) = \mathbf{A}(\mathbf{r}, t) + \nabla f(\mathbf{r}, t)$

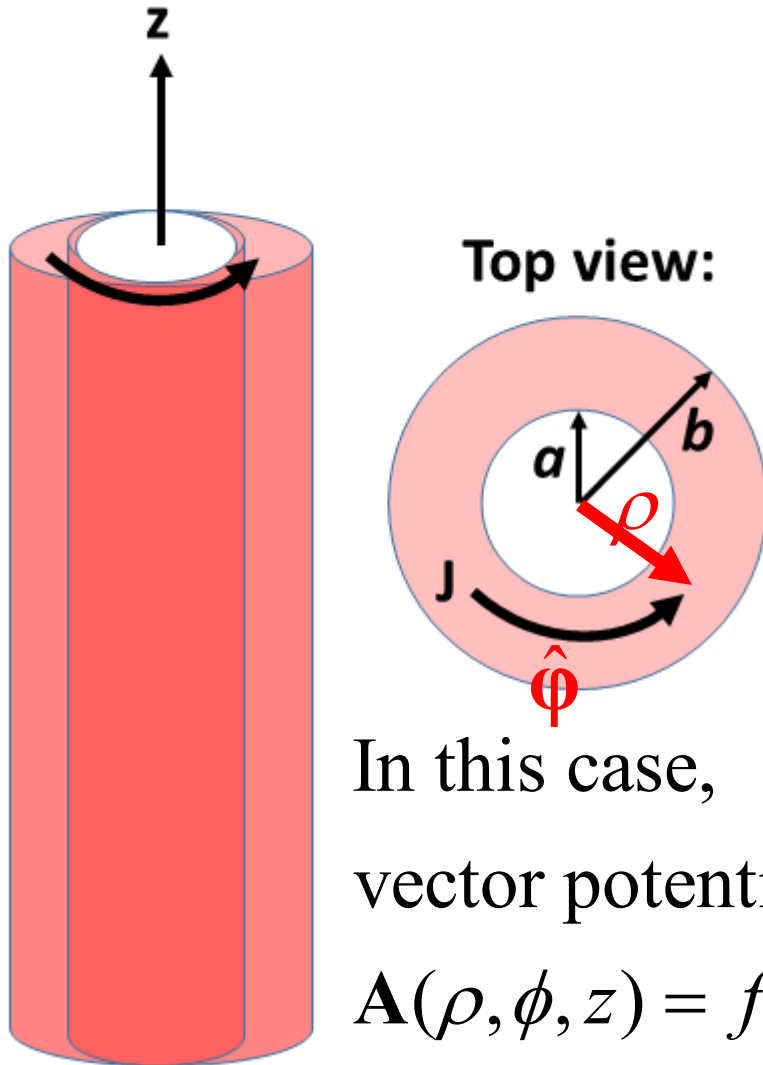
Note that $\nabla \times \mathbf{A}(\mathbf{r}, t) = \nabla \times \mathbf{A}'(\mathbf{r}, t)$

In our case, $\mathbf{A}(\mathbf{r}, t) = \frac{1}{2} B_0 (-y \hat{\mathbf{x}} + x \hat{\mathbf{y}})$

$$\mathbf{A}'(\mathbf{r}, t) = -B_0 y \hat{\mathbf{x}}$$

What is $f(\mathbf{r}, t)$?

To Julia's question about representing the electromagnetic terms in "generalized" coordinates



In this case, it is convenient to represent the vector potential in cylindrical coordinates:

$$\mathbf{A}(\rho, \phi, z) = f(\rho)\hat{\phi}$$

Now consider formulation of motion with constraints --
Comments on generalized coordinates:

$$L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} = 0$$

Here we have assumed that the generalized coordinates q_σ are independent. Now consider the possibility that the coordinates are related through constraint equations of the form:

Lagrangian : $L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$

Constraints : $f_j = f_j(\{q_\sigma(t)\}, t) = 0$

Modified Euler - Lagrange equations : $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} + \sum_j \lambda_j \frac{\partial f_j}{\partial q_\sigma} = 0$

Lagrange
multipliers



Some details --

Lagrangian : $L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$

Constraints : $f_j = f_j(\{q_\sigma(t)\}, t) = 0$

Modified Euler - Lagrange equations : $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} + \sum_j \lambda_j \frac{\partial f_j}{\partial q_\sigma} = 0$

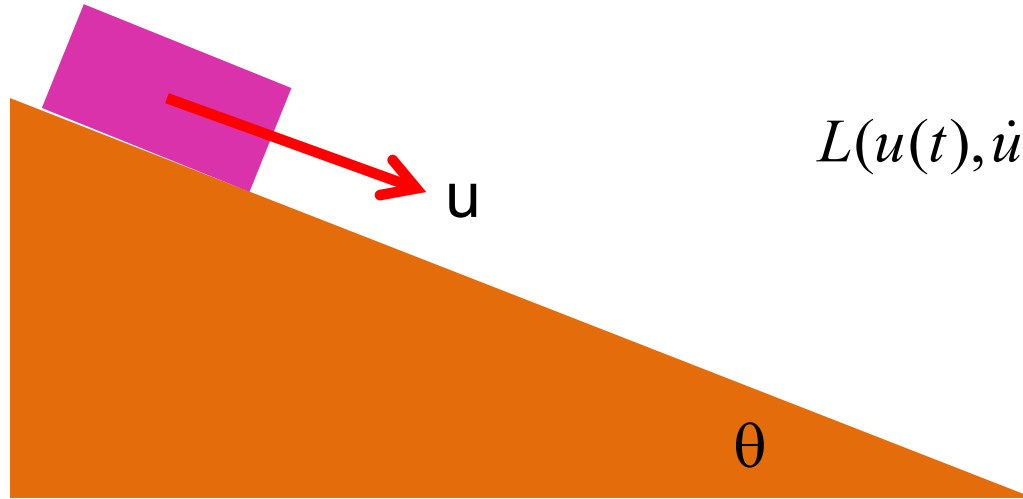
This amounts to modifying our optimization problem --

$\delta S = 0$ and for each i : $\delta f_i = 0$

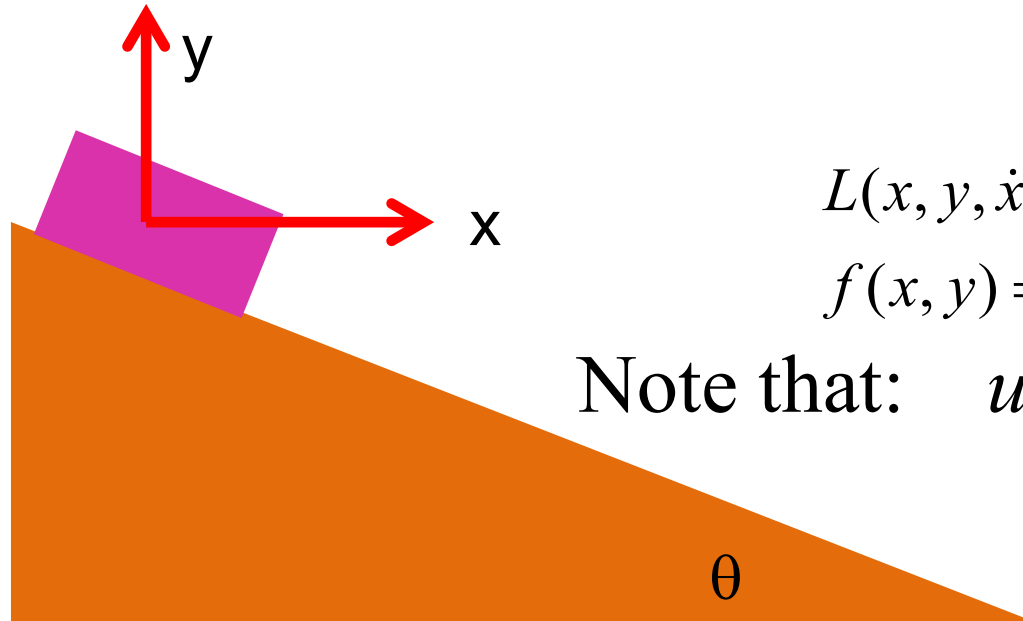
$\Rightarrow \delta W \equiv \delta(S + \sum_i \lambda_i f_i) = 0$, introducing the new constants λ_i .



Simple example:



$$L(u(t), \dot{u}(t)) = \frac{1}{2} m \dot{u}^2 + m g u \sin \theta$$



$$L(x, y, \dot{x}, \dot{y}) = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - m g y$$

$$f(x, y) = \sin \theta x + \cos \theta y = 0$$

Note that: $u = x \cos \theta - y \sin \theta$

Case 1:

$$L(u(t), \dot{u}(t)) = \frac{1}{2} m \dot{u}^2 + m g u \sin \theta$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{u}} - \frac{\partial L}{\partial u} = 0 = m \ddot{u} - m g \sin \theta = 0$$

$$\Rightarrow \ddot{u} = g \sin \theta$$

Case 2:

$$L(x, y, \dot{x}, \dot{y}) = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - m g y$$

$$f(x, y) = \sin \theta x + \cos \theta y = 0$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} + \lambda \frac{\partial f}{\partial x} = 0 = m \ddot{x} + \lambda \sin \theta$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} + \lambda \frac{\partial f}{\partial y} = 0 = m \ddot{y} + m g + \lambda \cos \theta$$

$$\sin \theta \ddot{x} + \cos \theta \ddot{y} = 0$$

$$\Rightarrow \lambda = -m g \cos \theta$$

$$(\cos \theta \ddot{x} - \sin \theta \ddot{y}) = g \sin \theta$$

Which method would you use to solve the problem?

Case 1

Case 2

Force of constraint;
normal to incline

Rational for Lagrange multipliers

Recall Hamilton's principle:

$$S = \int_{t_i}^{t_f} L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t) dt$$

$$\delta S = 0 = \int_{t_i}^{t_f} \left(\sum_{\sigma} \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} \right) \delta q_\sigma \right) dt$$

With constraints: $f_j = f_j(\{q_\sigma(t)\}, t) = 0$

Variations δq_σ are no longer independent.

$$\delta f_j = 0 = \sum_{\sigma} \frac{\partial f_j}{\partial q_\sigma} \delta q_\sigma \quad \text{at each } t$$

\Rightarrow Add 0 to Euler-Lagrange equations in the form:

$$\sum_j \lambda_j \sum_{\sigma} \frac{\partial f_j}{\partial q_\sigma} \delta q_\sigma$$

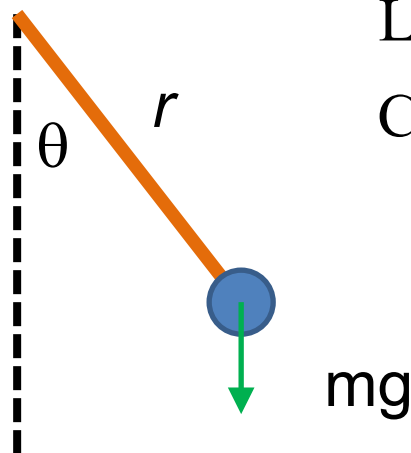
Euler-Lagrange equations with constraints:

Lagrangian: $L = L(\{q_\sigma(t)\}, \{\dot{q}_\sigma(t)\}, t)$

Constraints: $f_j = f_j(\{q_\sigma(t)\}, t) = 0$

Modified Euler - Lagrange equations: $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\sigma} - \frac{\partial L}{\partial q_\sigma} + \sum_j \lambda_j \frac{\partial f_j}{\partial q_\sigma} = 0$

Example:



Lagrangian: $L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) + mgr \cos \theta$

Constraints: $f = r - \ell = 0$

Example continued:

$$\text{Lagrangian: } L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) + mgr \cos \theta$$

$$\text{Constraints: } f = r - \ell = 0$$

$$\frac{d}{dt} m \dot{r} - m r \dot{\theta}^2 - mg \cos \theta + \lambda = 0$$

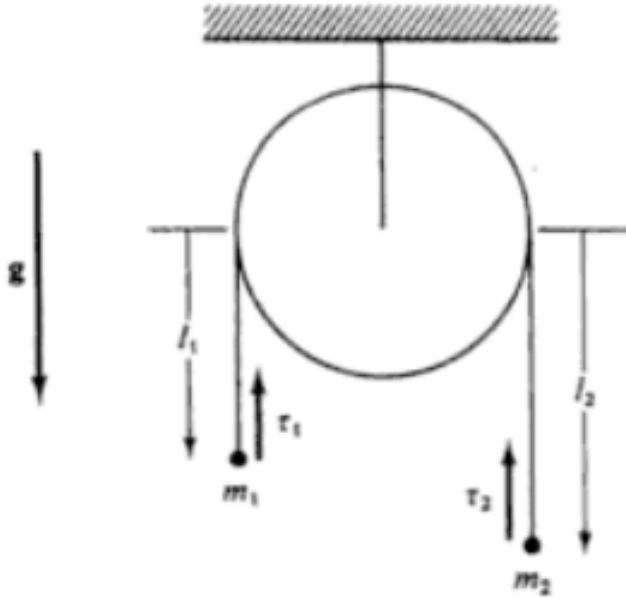
$$\frac{d}{dt} m r^2 \dot{\theta} + mgr \sin \theta = 0$$

$$\dot{r} = 0 = \ddot{r} \quad r = \ell$$

$$\Rightarrow \ddot{\theta} = -\frac{g}{\ell} \sin \theta$$

$$\Rightarrow \lambda = m \ell \dot{\theta}^2 + mg \cos \theta$$

Another example:



Lagrangian: $L = \frac{1}{2} m_1 \dot{\ell}_1^2 + \frac{1}{2} m_2 \dot{\ell}_2^2 + m_1 g \ell_1 + m_2 g \ell_2$

Constraints: $f = \ell_1 + \ell_2 - \ell = 0$

$$\frac{d}{dt} m_1 \dot{\ell}_1 - m_1 g + \lambda = 0$$

$$\frac{d}{dt} m_2 \dot{\ell}_2 - m_2 g + \lambda = 0$$

$$\dot{\ell}_1 + \dot{\ell}_2 = 0 = \ddot{\ell}_1 + \ddot{\ell}_2$$

Figure 19.1 Atwood's machine.

$$\Rightarrow \lambda = \frac{2m_1 m_2}{m_1 + m_2} g$$

$$\ddot{\ell}_1 = -\ddot{\ell}_2 = \frac{m_1 - m_2}{m_1 + m_2} g$$